An *Ab Initio* Investigation of Nickel(II) Complexes with Water, Chloride, Hydroxide and Ammonia

By Liwei Cheng

Abstract

One of the main undesirable processes in Supercritical Water-Cooled Reactors (SCWR) is the possible formation of metal ligand complexes from ions present in coolant which is very poorly understood both experimentally and theoretically. As pressurized and heated solutions are difficult to work with, computational methods have now become an important research tool in this respect. A series of *ab initio* calculations of nickel(II) with hydroxide, ammonia, water and chloride have been performed at HF, MP2 and B3LYP levels with 6-31+G* basis sets. A thorough examination of geometries, energies and vibrational frequencies has been carried out for all species up to and including hexacoordinate species. The computational results are compared with experimental data where available.

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<tr>
<td>ACEnet</td>
<td>Atlantic Computational Excellence Network</td>
</tr>
<tr>
<td>AECL</td>
<td>Atomic Energy of Canada Limited</td>
</tr>
<tr>
<td>B3LYP</td>
<td>Becke 3-Parameter Lee-Yang-Parr Hybrid Functional</td>
</tr>
<tr>
<td>CANDU</td>
<td>Canada Deuterium Uranium</td>
</tr>
<tr>
<td>COSMO</td>
<td>Conductor-like Screening Model</td>
</tr>
<tr>
<td>C-PCM</td>
<td>Conductor-like Polarizable Continuum Model</td>
</tr>
<tr>
<td>DFT</td>
<td>Density Functional Theory</td>
</tr>
<tr>
<td>GTO</td>
<td>Gaussian-Type Orbital</td>
</tr>
<tr>
<td>HF</td>
<td>Hartree-Fock</td>
</tr>
<tr>
<td>HPC</td>
<td>High-Performance Computing</td>
</tr>
<tr>
<td>LCAO</td>
<td>Linear Combination of Atomic Orbitals</td>
</tr>
<tr>
<td>LWR</td>
<td>Light Water Reactor</td>
</tr>
<tr>
<td>MO</td>
<td>Molecular Orbital</td>
</tr>
<tr>
<td>MP2</td>
<td>Second-order Møller-Plesset Perturbation Theory</td>
</tr>
<tr>
<td>NMR</td>
<td>Nuclear Magnetic Resonance</td>
</tr>
<tr>
<td>PCM</td>
<td>Polarizable Continuum Model</td>
</tr>
<tr>
<td>SCF</td>
<td>Self-Consistent Field</td>
</tr>
<tr>
<td>SCWR</td>
<td>Supercritical Water-cooled Reactor</td>
</tr>
<tr>
<td>STO</td>
<td>Slater-Type Orbital</td>
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<tr>
<td>UV/Vis</td>
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Chapter 1: Introduction

1.1 Supercritical Water-Cooled Reactor

CANDU Supercritical Water-Cooled Reactors (SCWR) are one of the six innovative concepts selected for Gen IV nuclear reactors based on safety and environmental issues which are being developed by Atomic Energy of Canada Limited (AECL) [1]. The SCWR can be considered a hybrid between a LWR (Light Water Reactor) and a fossil-fuel SCW power plant. The power cycle proposed for SCWR would be a direct power cycle using a single phase coolant [2]. The SCWR concept is being built for several applications, such as hydrogen production, steam applications such as extraction of oil from oil sands, desalination, etc. [3] The SCWR offers many advantages such as the elimination of components like steam generators, steam separators, dryers, and a low coolant mass inventory resulting in smaller components with much higher thermal efficiency (about 45% versus about 33% efficiency for LWRs) [4]. One of the primary objectives is to lower unit energy costs via advancements in thermodynamic efficiency which is attained by increasing the coolant core mean temperature in excess of 500 °C, leading to plant thermodynamic efficiencies as high as 50%. Thus, the SCWR power cycle would be a direct cycle operating at a pressure of 25 MPa, with core inlet and outlet temperatures of about 550 K and as high as 900 K respectively [5].

Unfortunately, accompanying the increased efficiency, supercritical water at high operating temperatures and pressures is very reactive especially in an oxidizing environment. Under extreme conditions, the properties of water affected include the density, hydrogen bonding, ionization product, dielectric constant, heat capacity, and
transport properties. These properties have a large impact on the corrosion properties of metal alloys in the construction materials, resulting in anodic dissolution of the alloys. [6] In addition, certain aggressive ions (e.g. Cl⁻, Br⁻) can induce passivity breakdown, resulting in various forms of localized corrosion, such as pitting attack and stress corrosion cracking. [7] The transition metal in the metal alloys could be leached into the environment and form aqua ions, which combine with the surrounding anions such as Cl⁻ and OH⁻. However, there will be no phase separation in the SCWR coolant, and the corrosion product formed in the reactor core will be carried out of core with the supercritical coolant. Those corrosion products may lead to deposition in the fuel or turbines. [8]

1.2 Nickel

Nickel is a hard, silvery white metallic element, (relative atomic mass 58.69) found in the first transition series group VIIIb of the periodic table. Nickel was first isolated in 1751 in impure form by Axel F. Cronstedt, a Swedish chemist. [9] The density of nickel is 8.908 g/cm³, its melting point is 1453 ºC, and its boiling point is 2732 ºC. [10] Five natural isotopes are known, of which ⁵⁸Ni (68.27%) and ⁶⁰Ni (26.10%) are the most abundant. The common oxidation states of nickel are 0 and +2, but other states (+1, +3, +4, and +6) are found in a few compounds. However, the +2 oxidization state is the most prevalent form in aqueous solution. Nickel(VI) is known in NiO₄²⁻, which is a very powerful oxidizing agent and will oxidize water. Nickel (IV) is found in NiO₂, which is also a very strong oxidizing agent. Some nickel(III) species, such as NiO(OH), are also
known and are strong oxidizing agents. Besides these species, various oxidization states of nickel are found in complex ions. [11]

Nickel-based alloys are vitally important to modern industry due to their ability to withstand a wide variety of severe operating conditions involving corrosive environments, high temperatures, high stress, and combinations of these factors. [12] In the nuclear industry, nickel-based alloys are chosen and used in gas turbine engines and a spectrum of other applications of relevant use. [13] In particular, nickel-based alloys are chosen because of their generally higher creep strength compared to austenitic and ferritic-martenstitic steels for the supercritical water reactor, [14] nickel-base alloys have been proposed for Gen IV reactor in-core components. [15]

The excellent corrosion resistance of nickel-based alloys is largely due to the surface protection in aqueous electrolytes or moist air by ultrathin dense oxide/hydroxide films. [16] However, the passive film does not afford complete protection of the substrate alloys. For instance, tribological contact or aggressive ions can often damage or even remove the passive film, and expose the active substrate alloy surface to the corrosion environment. [17] In this case, nickel(II) ion could be leached into the environment and form aqua complexes with surrounding anions and/or ammonia, which could be introduced by pH adjustment. The anions could be introduced through the cooling system (i.e. Cl from river water) or from the addition of water itself (OH ions). These nickel complexes can be transported out of core, and may deposit in the fuel or turbines and corroding these areas. Therefore, it is necessary to carry out the investigation of the possible complexes nickel will form in SCWR for both scientific research and
engineering practices. As pressurized and heated solutions are difficult to work with, computational methods have now become an important research tool.

1.3 Computational Chemistry

1.3.1 Quantum Mechanics

Computational chemistry is based on quantum mechanics, which provides a very detailed picture of molecules. The postulates and theorems of quantum mechanics form the rigorous foundation for the prediction of observable chemical properties from first principles. The fundamental postulates of quantum mechanics assert that microscopic systems are described by wave functions that completely characterize all of the physical properties of the system. [26]

In the late seventeenth century, Isaac Newton summarized classical mechanics, the laws of motion of macroscopic objects. In the early twentieth century, physicists found that classical mechanics does not correctly describe the behavior of very small particles such as the electrons and nuclei of atoms and molecules. [18] Quantum mechanics determines the properties of nanomaterials (objects with at least one dimension in the range 1 to 100 nm), and calculation methods to deal with nanomaterials are being developed. [18]

The development of quantum mechanics began in 1900 with Planck’s study of blackbody radiation, which is the light emitted by a heated solid that absorbs all light falling on it. [18] When physicists used classical mechanics to predict the intensity-
versus-frequency curve for emitted blackbody radiation, they found a result in complete
disagreement with the high-frequency portion of the experimental curves. [18] Planck
assumed the radiation emitters and absorbers in the blackbody to be harmonically
oscillating electric charges (resonators), and the total energy of those resonators whose
frequency is \( \nu \), consisted of \( N \) indivisible “energy elements” which led to curves that
agreed with the experimental blackbody curves. [18] Thus the energy of each resonator
was quantized, meaning that only certain discrete values were allowed for resonator
energy. [18]

The second application of energy quantization was to the photoelectric effect, in
which light shining on a metal causes emission of electrons. [18] Physicists found it
difficult to reconcile their observations with the classical electromagnetic wave theory of
light. An explanation was suggested in 1905 by Einstein who proposed that the incident
light be viewed as being comprised of discrete units of energy. [19] The photoelectric
effect shows that light can exhibit particlelike behavior in addition to the wavelike
behavior it shows in diffraction experiments. [18] De Broglie proposed that all material
particles are associated with waves, but that the existence of these waves is likely to be
observable only in the behavior of extremely light particles. [19] Quantum mechanics is
required for those microscopic particles as classical mechanics only applies to
macroscopic particles.

Microscopic particles have a wave nature. When such a particle is located by
using radiation of photons of large momentum, the interaction of photon and the electron
changes the momentum of the electron. [20] Thus both position and velocity of the
particle cannot be measured simultaneously. Heisenberg [21] stated the uncertainty principle, which can be derived from quantum theory,

$$\Delta p \Delta x \geq \frac{\hbar}{2}$$ (1)

That means to locate an electron to within a region $\Delta x$, there will be an uncertainty in the momentum $\Delta p$ of the electron. [22]

Therefore, a wave equation is needed for quantum systems. Schrödinger [23] obtained such an equation which is written as:

$$\hat{H}\Psi = E\Psi$$ (2)

In the Schrödinger equation, $\Psi$ is used for an unspecified quantum mechanical wave function, $E$ is the energy eigenvalue. $\hat{H}$ is called the Hamiltonian operator. The wave function $\Psi$ is an eigenfunction of the Hamiltonian operator, and the energy $E$ is an eigenvalue of the Hamiltonian operator as the result of operating $\hat{H}$ on $\Psi$ is simply to give $\Psi$ back again, only multiplied by a constant factor, the energy $E$. The Hamiltonian operator is written as:

$$\hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x)$$ (3)

In this equation, the $-\frac{\hbar^2}{2m} \frac{d^2}{dx^2}$ is the kinetic energy part, and $V(x)$ is the potential energy part. $m$ is the mass of the particle at position $x$.

Equation 2 does not contain time and is called the time-independent Schrödinger equation. The wave function obtained from equation 2 is called the stationary-state wave
function, and many problems of chemical interest can be described in terms of stationary-
state wave functions. [21] The wave function contains all possible information of a
system. Also, the one-dimensional, one-particle systems can be extended to three-
dimensional, many-particle systems. For a three-dimensional, one-particle system, the
Schrödinger equation is:

\[- \frac{\hbar^2}{2m} \left( \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} \right) + V(x, y, z) = E\psi(x, y, z) \quad (4)\]

The first term of the Hamiltonian operator in equation 4 is called Laplacian
operator, denoted as \(\nabla^2\):

\[\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \quad (5)\]

The Schrödinger equation for many-particle, three-dimensional systems are to
sum up all the kinetic energy and potential energy components:

\[\left[ - \sum_{i=1}^{n} \frac{\hbar^2}{2m_i} \nabla_i^2 + V(x_1, y_1, z_1, ..., x_i, y_i, z_i) \right] = E\psi \quad (6)\]

The quantum mechanical description of a hydrogen atom is the starting point for
computational chemistry, and it can be solved analytically. [24] However, atoms with
more than one electron cannot be solved exactly as there is a interelectron repulsion term
in the Schrödinger equation. In order to simplify the problem, the Born-Oppenheimer
approximation is invoked. [25] As the nuclei of molecular systems are moving much
more slowly than the electrons, it is convenient to decouple the motions of nuclei, and
compute electronic energies for fixed nuclear position. [26] Therefore, to calculate the
energy of a molecule is to solve the electronic Schrödinger equation and then add the electronic energy to the internuclear repulsion to get the total internal energy. [27] A deeper consequence of the Born-Oppenheimer approximation is that a molecule has a shape. [27]

1.3.2 Hartree-Fock Theory

The Hartree-Fock (HF) calculation is the simplest kind of ab initio calculation which was first performed on atoms by Hartree in 1928. [28] The fundamental assumption of HF theory, that each electron sees all of the others as an average field, allows for tremendous progress to be made in carrying out practical molecular orbital calculations. [26] The HF method is used to solve Schrödinger equations of atoms and molecules and predict their properties such as geometries and vibrational frequencies. According to the Pauli Exclusion Principle, the wave function for a HF calculation is one or more antisymmetrized products of one-electron spin-orbitals. A convenient way to produce an antisymmetrized product is to use Slater determinants:

$$
\Psi(1, 2, \ldots, 2N) = \frac{1}{\sqrt{(2N)!}} \begin{vmatrix} 
\psi_1(1) & \psi_1(1) & \cdots & \psi_N(1) & \psi_N(1) \\
\psi_1(2) & \psi_1(2) & \cdots & \psi_N(2) & \psi_N(2) \\
\vdots & \vdots & \ddots & \vdots & \vdots \\
\psi_1(2N) & \psi_1(2N) & \cdots & \psi_N(2N) & \psi_N(2N) 
\end{vmatrix}
$$

(7)

2N is the number of electrons in the closed shell system as the wave functions are represented by N doubly occupied spatial orbitals. \(\alpha\) and \(\beta\) are spin orbitals, and \(\frac{1}{\sqrt{(2N)!}}\) is the normalization constant. The wavefunctions in the Slater determinants are functions that are properly antisymmetric with respect to exchange of coordinates of any pair of
electrons. Such an exchange would correspond to interchanging two rows of the Slater determinant according to property of determinants that their value changes sign if two rows are exchanged. [24] If there are two identical spin-orbitals, then there will be two identical columns in the Slater determinant according to the property of determinants that their value is zero if two columns are identical. [24] It follows from the Schrödinger equation that the energy of a system is given by:

\[ E = \int \frac{\psi^* \hat{H} \psi}{\psi^* \psi} \, dt. \]  

(8)

Using Dirac notation it can be written:

\[ E = \langle \Psi(1, 2, \ldots, 2N) | \hat{H} | \Psi(1, 2, \ldots, 2N) \rangle, \]

(9)

while the Hamiltonian operator for a molecule with 2N electrons and M atomic nuclei (the \( M^{th} \) nucleus has the charge \( Z_M \)) can be written as:

\[ \hat{H} = \sum_{i=1}^{2N} \left( -\frac{1}{2} \nabla_i^2 + \sum_{l=1}^{M} \frac{Z_M}{r_{Mi}} + \sum_{l=1}^{N} \sum_{j=l}^{N} \frac{1}{r_{ij}} \right). \]

(10)

The Hamiltonian is composed of electron kinetic energy terms, nucleus-electron attraction potential energy terms, and electron-electron repulsion potential energy terms. From the Born-Oppenheimer approximation, the nucleus-nucleus repulsion potential energy terms have been omitted, and these can simply be added to the electronic energy as they are constants. The energy can be written as:

\[ E = 2 \sum_{l=1}^{N} H_{ll} + \sum_{l=1}^{N} \sum_{j=1}^{N} (2J_{lj} - K_{ij}). \]

(11)
The terms in Eq. 11 have these meanings

\begin{align}
H_{ii} &= \int \psi_i^*(1) \left( -\frac{1}{2} \mathbf{\nabla}_1^2 - \sum_{a,ll} \frac{Z_a}{r_{1a}} \right) \psi_i(1) dr \\
J_{ij} &= \int \psi_i^*(1) \psi_i(1) \left( \frac{1}{r_{12}} \right) \psi_j(2) \psi_j^*(2) dr_1 dr_2 \\
K_{ij} &= \int \psi_i^*(1) \psi_j^*(2) \left( \frac{1}{r_{12}} \right) \psi_i(2) \psi_j(1) dr_1 dr_2
\end{align}

\( H_{ii} \) is the electronic energy of a single electron moving simply under the attraction of a nuclear core, with all the other electrons stripped away. \( J_{ij} \) is called a Coulomb integral, and it represents the electrostatic repulsion between an electron in \( \psi_i \) and one in \( \psi_j \). \( K_{ij} \) is called an exchange integral, and it differs only in exchange of electrons.

The energy calculated from Eq. 9 will be the exact, true energy of the molecule only if the wavefunction and the Hamiltonian are exact. The variation principle states that the energy calculated from Eq. 9 must be greater than or equal to the true ground-state energy of the molecule. [27] The energy associated with the wave function is determined from all the coefficients. The Hartree-Fock method follows a SCF (self-consistent field) procedure, where the orbital coefficients are first estimated and then iterated to convergence.

### 1.3.3 Basis Sets

The basis set is the set of mathematical functions (basis functions), linear combinations of which yield molecular orbitals. [27] The functions are usually centred on
nuclei. Approximating molecular orbitals as linear combinations of basis functions is called the LCAO-MO approach. [27] However, the functions are not necessarily conventional atomic orbitals, they can be any set of mathematical functions that are convenient to manipulate and which in linear combination give useful representations of molecular orbitals. Physically, several basis functions describe the electron distribution around an atom and combining atomic basis functions yields the electron distribution in the molecule as a whole. [27] The electron distribution around an atom can be represented in two ways, Slater functions and Gaussian functions. [29] Slater functions are used in semiempirical methods, and modern molecular ab initio programs employ Gaussian functions.

The essence of STO (Slater-type-orbital) is to place on each nucleus one or more STOs. [27] Generally, the larger the number of STOs and/or the greater the care taken in selecting orbital exponents, the more accurate the final wavefunction and the energy will be. The mathematical form of a normalized STO is:

\[
S(r, \theta, \zeta, \phi, n, l, m) = N_{nl} r^{n-1} e^{-\zeta r} Y_l^m(\theta, \phi),
\]

where \( \zeta \) is an exponent that can be chosen according to a simple set of rules developed by Slater that depend on the atomic number. [30] \( N_{nl} \) is the normalization constant. \( r \) is the distance between the nucleus and the electron. \( n \) is the principal quantum number for the valence orbital, and the spherical harmonic functions \( Y_l^m(\theta, \phi) \) depend on the angular momentum quantum numbers \( l \) and \( m \).

In ab initio HF theory, STOs suffer from a fairly significant limitation. There is no analytical solution available for the general two-electron integral. [26] The requirement
that such integrals be solved by numerical methods severely limits their utility in molecular systems of any significant size. To speed up molecular integral evaluation, Boys [31] proposed in 1950 the use of Gaussian-type functions instead of STOs for the atomic orbitals in an LCAO wave function. The Gaussian-type functions can be written as:

\[ G(r, \theta, \alpha, \phi, n, l, m) = N_l r^l e^{-\alpha r^2} Y_l^m(\theta, \phi), \]  

(16)

where \( \alpha \) is the shielding constant which is an exponent controlling the width of the Gaussian type orbital (GTO). In GTO, the radial decay of the STOs are changed from \( e^{-r} \) to \( e^{-r^2} \). Gaussian integral evaluation takes much less computer time than Slater integral evaluation because the product of two Gaussians on two centres is a Gaussian on a third centre. [27] Thus all three- and four- centre two electron repulsion integrals are reduced to two-centre integrals. However, there are two obvious problems connected with using Gaussian functions as basis functions: one is that they do not have cusps at \( r=0 \) as s-type hydrogen-like atomic orbitals (AOs) do, the other is that they decay faster at larger \( r \) than do hydrogen-like AOs. [19] This led to a practice of replacing each STO in a basis set by a number of Gaussian functions. When a basis function is defined as a linear combination of Gaussians, it is referred to as a contracted basis function, and the individual Gaussians from which it is formed are called primitive Gaussians.

There are many terms used to describe a STO basis set. A minimal basis set consists of one STO for each inner-shell and valence-shell AO of each atom. [18] During most molecular bonding, it is the valence electrons which principally take part in the bonding. In recognition of this fact, a split-valence basis set uses two or more STOs for
each valence AO but only one STO for each inner-shell (core) AO. Upon molecule formation, AOs are distorted in shape and have their centres of charge shifted, in asymmetric molecule environments orbitals become distorted from their atomic shapes (polarization). A polarized basis set is adding basis function STOs whose angular quantum numbers are greater than the maximum angular quantum number of the valence shell of the ground-state atom, which is denoted by “*”. [18] Anions tend to have their orbitals expanded compared to neutral molecules, so diffuse functions need to be included. A basis set with diffuse Gaussian functions with fairly small orbital exponents is designated by “+”. [22]

1.3.4 Møller–Plesset perturbation theory

The Hartree-Fock wave function takes into account the interactions between electrons only in an average way which does not consider the instantaneous interactions between electrons. Møller–Plesset perturbation theory is a post-Hartree-Fock method which treats correlated motion of electron pairs better than the Hartree-Fock method. The Møller–Plesset treatment of electron correlation is based on perturbation theory, and this particular approach was described by Møller and Plesset in 1934 [32] and developed into a practical molecular computational method by Binkley and Pople in 1975. [33] There is a hierarchy of Møller–Plesset energy levels: MP0, MP1, MP2, etc… which successively account more thoroughly for interelectronic repulsion. [27] MP0 would use the electronic energy obtained by simply summing the Hartree-Fock one-electron energies, and MP1 is just the Hartree-Fock energy. [27] MP2 is the first MP level to go beyond the Hartree-
Fock treatment, and the MP2 energy is the Hartree-Fock energy plus a correction term that represents a lowering of energy brought about by allowing the electrons to avoid one another better than in Hartree-Fock treatment:

\[ E_{MP2} = E_{HF} + E^{(2)} \]  \hspace{1cm} (17)

The \( E^{(2)} \) is a purely electronic term which represents the perturbation correction, and it can be written as:

\[ E^{(2)} = \sum_{s \neq 0} \frac{|\langle \psi_s^{(0)} | \hat{H} | \phi_0 \rangle|^2}{E_0 - E_s^{(0)}}, \]  \hspace{1cm} (18)

where \( \phi_0 \) is the ground-state Hartree-Fock wave function, \( E_0 \) is the zeroth order MP energy, \( \psi_s^{(0)} \) is the unperturbed wave function, and \( E_s^{(0)} \) is the correlated energy. The perturbation \( \hat{H}' \) is the difference between the true interelectronic repulsions and the Hartree-Fock interelectronic potential.

Calculated properties like geometries and relative energies tend to be better when done with Møller–Plesset perturbation theory. [27] The essence of the Møller–Plesset method is that the correction term handles electron correlation by promoting electrons from occupied to unoccupied MOs, giving electrons more room to move and thus making it easier for them to avoid one another, and the decreased interelectronic repulsion results in a lower electronic energy. [27]

MP3 calculations take a lot longer than MP2 calculations, but provide little improvement over MP2 molecular properties and so are rarely done. By far, the most common MP level used is MP2. [18]
1.3.5 Density Functional Theory

Density functional theory is based on the two Hohenberg-Kohn [34] theorems, which state that the ground-state properties of an atom or molecule are determined by its electron density function, and that a trial electron density must give energy greater than or equal to the true energy. [27] Density functional theory is based on the electron probability density or the charge density, and designated by $\rho(x, y, z)$. [27] This electron density $\rho$ is measureable by X-ray diffraction or electron diffraction. [35] The electron density is a function of position only, that is, no matter how big the molecule may be, the electron density remains a function of three variables (while the wavefunction of an n-electron molecule will have $4n$ variables). Thus, the density trumps the wavefunction in three ways: it is measurable; it is intuitively comprehensible, and it is mathematically more tractable. [27] The main advantage of density functional theory is that in about the same time needed for a Hartree-Fock calculation one can often obtain results of about the same quality as from MP2 calculations. [27]

The ground state electronic energy $E_o$ is a functional of $\rho_o$ and is written as:

$$E_o = E_o[\rho_o], \quad (19)$$

where the square brackets denote a functional relation. The ground state electronic energy of the real molecule is the sum of the electron kinetic energies, the nucleus-electron attraction potential energies, and the electron-electron repulsion potential energies:

$$E_o = \langle T[\rho_o] \rangle + \langle V_{Ne}[\rho_o] \rangle + \langle V_{ee}[\rho_o] \rangle \quad (20)$$
The angle brackets indicate that these energy terms are quantum-mechanical average values, and each is a functional of the ground-state electron density. For an accurate treatment of the electronic kinetic energy term, Kohn and Sham [36] proposed the density in terms of one-electron orbitals $\phi$. Then the electron density can be written as:

$$\rho_s = \sum_i^n |\phi_i|^2$$  \hspace{1cm} (21)

The electron-electron repulsion term can be decomposed into Coulomb and exchange terms, $J[\rho]$ and $K[\rho]$. The final DFT energy expression then can be written as:

$$E_{\text{DFT}} = T_s[\rho] + V_{\text{Ne}}[\rho] + J[\rho] + E_{\text{xc}}[\rho] ,$$ \hspace{1cm} (22)

where the exchange correlation functional $E_{\text{xc}}[\rho]$ contains the difference between the exact kinetic energy and $T_s$, the exchange part of electron-electron repulsions, $K[\rho]$, and correlations to both $K[\rho]$ and $J[\rho]$.

In the Kohn-Sham case, the HF exchange operators are replaced by the functional derivative of the exchange correlation energy.[19] The exact form of $E_{\text{xc}}[\rho]$ is not currently known, but a rapidly growing list of approximate exchange correlation functionals have appeared in the literature. [19] Generally, most existing exchange correlation functionals are split into a pure exchange and correlation contribution, $E_x[\rho]$ and $E_c[\rho]$. The functional used in this research is B3LYP method which is a hybrid functional. B3LYP stands for Becke, three-parameter, Lee-Yang-Parr, and it includes 20% exact exchange and involves three semiempirical parameters that were obtained by fits to experimental thermochemical data of small molecules.[37]
1.3.6 **Conductor-like Screening Model (COSMO)**

All the previous methods of calculations have been done in the gas phase. However, the properties of molecules and transition states can differ considerably between the gas phase and solution. [38] COSMO is a calculation method for determining the electrostatic interaction of a molecule with a solvent which was developed by Klamt and Schüürmann.[39]

The COSMO model is a dielectric model in which the solute molecule is embedded in a molecule-shaped cavity surrounded by a dielectric medium with given dielectric constant $\varepsilon$. [27] The key algorithms for calculating the properties of a molecule in solution is to formulate a solution Hamiltonian operator in which these energy terms appear in addition to the in vacuo terms of the electron kinetic energy, electron-nucleus attraction, and electron-electron repulsion. [27] The COSMO model uses a conducting medium ($\varepsilon$ infinite) and introduces the solvent dielectric constant as a correction factor. [39] The COSMO method is very fast when used, as is usually the case, with gas phase geometries followed by single-point calculations in solvent. The method used in this research is a particular implementation of COSMO that allows efficient geometry optimization in solution is called C-PCM (conductor polarizable continuum model).
1.4 Objectives

The primary objective of this research is to perform a comprehensive ab initio computational study of nickel(II) complexes with water, chloride, hydroxide and ammonia. These series of complexes will be studied up to and including the hexacoordinate species. The studies will include geometry optimization and bond length and frequency analysis. Single point C-PCM calculations will be performed based on the optimized B3LYP/6-31+G* structures. Calculated Raman spectra will be plotted for the optimized HF/6-31+G* complexes. Although this research will not include data at high temperature and pressure, it will provide a good starting point for the experimental work that will be performed at the University of Guelph by the Peter Tremaine group. Even with the Raman data being calculated at zero Kelvin and in the gas phase it will still be applicable to the high temperature, high pressure experiments because the frequencies will not actually deviate that much due to anharmonicity and weaker molecular bonds at high temperature.
Chapter II: Methods

All the calculations were performed using the Gaussian 03 software package on Atlantic Computational Excellence Network (ACEnet). [40]

Gaussian is a computational chemistry software program initially released in 1970 by John Pople and his research group as Gaussian 70 [41], and it has been continuously updated since then. The Gaussian package can do a variety of calculations, like molecular mechanics, semi-empirical calculations, ab initio calculations, etc. In this research, the Unix/Linux version of Gaussian 03 was used. ACEnet is a consortium of Atlantic Canada Universities providing researchers with high performance computing resources. [42] ACEnet has six primary clusters in Atlantic Canada, and the specific cluster used in this research is located at Memorial University of Newfoundland and is named “Placentia”.

The levels of theory used in this research are Hartree-Fock (HF), second order Møller–Plesset perturbation theory (MP2) and B3LYP hybrid density functional theory, and the basis set is 6-31+G* which is widely used for studies of transition metal containing systems. [120] Geometries were optimized using a stepping stone approach, in which the geometries at HF, MP2 and B3LYP levels were optimized sequentially. When this approach failed (as one or more ligands dissociated), the problematic level was skipped. Default optimization specifications were normally used. After each level, a frequency calculation was performed at the same level and the resulting Hessian was used in the following optimization (geom=allcheck). The frequencies depend on the second derivative of the energy with respect to the nuclear positions. A minimum on the
potential energy surface has all the force constants positive. [27] Since a frequency calculation involves taking square root of a force constant, and the square root of a negative number is an imaginary number, and it is called imaginary frequency. [38] Z-matrix coordinates constrained to the appropriate symmetry were used as required to speed up the optimizations. The calculated potential energy surface stationary points were characterized by the vibrational frequencies from the Hessian. All frequencies were real for a true minimum. Optimizations were carried out to completion for all levels in the sequence even if one or more of the levels proved not to be local minima. The structures were visualized with the MOLDEN software package. [43]

The polarized Raman spectra are plotted using a code created by my supervisor, Dr. Cory C. Pye. These spectra were created from the frequency data obtained in the output file of the HF/6-31+G* calculation primarily because Raman intensity calculations are standard in Hartree-Fock frequency calculations. Conductor-like polarizable continuum model (C-PCM) calculations were performed only on the stable, optimized B3LYP/6-31+G* structures. These calculations were single-point calculations and are completed by inserting SCRF=CPCM in the input line of the calculation along with the level of theory and basis set. Water is used as the default solvent.
Chapter III: Hydrated Nickel(II) and Nickel(II) Complexes Combines with Water, Chloride, Hydroxide and Ammonia

The ab initio investigation of nickel(II) complexes throughout this thesis involve four ligands and the complexes that they make with the nickel(II) ion. These ligands are hydroxide (OH\textsuperscript{-}), water (H\textsubscript{2}O), chloride (Cl\textsuperscript{-}) and ammonia (NH\textsubscript{3}). These ligands, and the complexes of them, are all important with respect to the development of the Supercritical Water-Cooled Reactor (SCWR). The formation of these complexes and their transportation throughout the reactor environment may lead to enhanced corrosion and deposition of the pipes and valves and, ultimately, a nuclear leakage.

This chapter includes the computational results and discussion/literature comparison of the nickel(II) complexes with water, chloride, hydroxide and ammonia. These aqua complexes were studied up to and including the hexa-coordinate species. The chloro, hydroxo, and ammine complexes were studied up to hexa-coordinate species with and without water molecules. The water molecules were added to investigate how hydration affects the stability of the given species. All of the results regarding total energies, bond lengths and vibrational frequencies are from the HF, MP2 and B3LYP calculations utilizing the 6-31+G\* basis set. These calculations are all zero Kelvin, gas-phase calculations. Molecular geometries are presented for the MP2/6-31+G\* and B3LYP/6-31+G\* calculations. The aqueous C-PCM results are all from single-point calculations using the fully optimized B3LYP/6-31+G\* calculations. All simulated Raman spectra are from the vibrational frequency data calculated from HF/6-31+G\* because Raman intensities are standard for HF calculations. All structures shown will
only be those of the absolute energy minima, unless other local minima were found that
had noticeably different geometries.

3.1 Water, Hydroxide, Chloride and Ammonia Ligands

Ab initio calculations were performed on the ligands involved with the nickel(II)
(with the triplet spin state) complexes for reasons of stability comparison. Table 3A.1,
found in the supplementary material section, gives the total energies of the ligands
with the given levels of theory (the basis set will be excluded in all energy tables because only
results utilizing the 6-31+G* basis set will be reported). Also included in Table 3A.1 is
the total free energy in solution, calculated using C-PCM and the optimized B3LYP
geometry.

Figure 3-1: Optimized MP2 and B3LYP geometries for the Ni(II) ion and all ligands
used in the calculations of nickel(II) complexes.
$3.2 \quad \textbf{Aquanickel(II) Complexes, } [\text{Ni(H}_2\text{O)}_n]^2^+ \text{, where } n=1-6$

The hydration structure of the Ni(II) ion has been well investigated using various experimental and computational methods. Hydrated nickel(II) complexes are known to form octahedral geometry species with 6-coordinated water molecules at room temperature, as well as tetrahedral complexes with 4-coordinated water molecules at high temperature. [79] Hoffmann et al. [59] showed that the structure about Ni$^{2+}$ changes from 6- to 4-coordinate upon increasing the temperature from 25 °C to 425 °C by XAFS measurement. The majority of the studies on hydrated nickel(II) complexes show that Ni(II) ion is coordinated by six H$_2$O molecules in the first coordination sphere. The $[\text{Ni(H}_2\text{O)}_6]^2^+$ ion has been studied by infrared photodissociation spectroscopy [80], X-ray absorption fine structure [81], X-ray absorption near-edge structure [50], molecular dynamics method (MD) [45,50,81,82], Monte Carlo method (MC) [51], ab initio [44,51,53,55,56,61,83], water exchange [47,57], NMR [46], Raman Spectroscopy [68,69], neutron diffraction [58], and X-ray crystallography [48]. The results for the monohydrate [44,45,51,60,67,84] and pentahydrate [44,53,80] have also been reported. Fewer studies have been completed on the dihydrate, trihydrate and tetrahydrate ions. [44]

$3.2.1 \quad \textbf{Results}$

Stable structures were found for all aquanickel(II) complexes up to and including the hexaaquanickel(II) ion. The total energies, including the aqueous C-PCM energy, for all aqua complexes studied are given in Table 3A.2 in the supplementary materials. The
lowest energy optimized MP2 and B3LYP structures of these complexes are given in Figure 3-2 along with their respective symmetries.

Figure 3-2: Optimized MP2 and B3LYP geometries for [Ni(H$_2$O)$_n$]$^{2+}$, where n=1-6. All structures are similar with the exception of the dihydrate which is D$_{2d}$ at MP2 and D$_{2h}$ at B3LYP. (grey=nickel, red=oxygen, white=hydrogen)
The monohydrate has a preferred symmetry dependent on the level of theory. For MP2 the most stable geometry is $C_{2v}$, but the $C_{2v}$ structure has an imaginary frequency at B3LYP level. The B3LYP structure was most stable with $C_s$ symmetry. For the dihydrate, $D_{2h}$ and $D_{2d}$ structures were tried first, and $D_{2h}$ structure has lowest energy at B3LYP level, whereas the $D_{2d}$ structure has lowest structure at MP2 level. The preferred structure for triaquanickel was $C_3$ at both levels. The tetrahydrate has $C_s$ symmetry and the pentahydrate has $C_{2v}$ symmetry. Finally, the hexahydrate has a preferred $T_h$ geometry at both levels.

Plots were constructed of the Ni-O bond lengths and vibrational stretching frequencies and are shown in Figure 3-3, Figure 3-4 and Figure 3-5 for HF, MP2 and B3LYP respectively. The simulated polarized Raman spectrum for the hexaaquanickel(II) species can be found in Figure 3-6.

### 3.2.2 Discussion/Literature Comparison

The structure found for the monohydrate complex is consistent with the computational study by Bustamante et al. [44]. The Ni-O bond length they found is 1.892 Å, which is slightly longer than the MP2 and B3LYP results presented here of 1.882 Å and 1.861 Å respectively. This difference can be attributed to the fact that Bustamante et al. used BPW91 theory with a larger basis set, 6-311+G(d,p). Iuchi et al. [45] found a value of about 1.900 Å from MD simulation for the Ni-O distance which is also slightly longer than the result from this research. Magnusson and Moriarty [60] found the Ni-O
distance to be 1.906 Å using MP2/6-311+G**. No experimental data was found regarding the monohydrate complex.

The bond lengths for the dihydrate are 1.902 Å and 1.876 Å for MP2 (D$_{2d}$) and B3LYP (D$_{2h}$) respectively. Bustamante et al. [44] reported the 1.867 Å for D$_{2h}$ structure which is in agreement with the B3LYP result presented here. Overall the Ni-O bond has increased slightly when compared to the monohydrate. No experimental data was found to compare these results to.

The most stable structure for the trihydrate complex proved to be C$_3$ for both levels of theory. The Ni-O bond lengths for these structures were found to be 1.942 Å and 1.928 Å for MP2 and B3LYP respectively. Bustamante [44] found Ni-O bond length of about 1.931 Å which is in agreement with the MP2 and B3LYP results. The bond lengths are longer when compared to the mono- and dihydrate values. The result can be attributed to there being an increased number of ligands which would have the effect of repulsion forcing the bonds to be longer. Again, no experimental data was available for comparison.

The tetrahydrate species has C$_s$ symmetry which is most stable. The average Ni-O bond lengths are 1.989 Å and 1.986 Å for MP2 and B3LYP respectively. Bustamante et al. [44] reported a bond length of 1.990 Å for the C$_4$ geometry. The C$_4$ structure has also been tested, and it is stable at both levels, but the energies are slightly higher than the C$_s$ structure. Since the symmetry is different, the results from this research cannot be compared to those of Ref. [44]. Other than this study, no others were found regarding molecular structure of tetrahydrate species.
The pentahydrate species that proved to be most stable has \( C_{2v} \) symmetry. The average Ni-O bond lengths are 2.032 Å and 2.036 Å for MP2 and B3LYP respectively. Bustamante et al. [44] reported a bond length of 2.047 Å for the same symmetry which is also slightly longer than our results. Mare et al. [53] found the Ni-O distance to be 2.017 Å – 2.052 Å from different basis sets. These values are in agreement with our results.

The hexahydrate species is the most reported of the nickel(II) aqua ion in the literature. The most stable geometry for this ion has been found to be \( T_h \). The Ni-O distances that have been found are 2.071 Å and 2.079 Å for MP2 and B3LYP respectively. There have been many studies on the hexaaquanickel(II) ion. Bustamante et al. [44] reported a bond length of 2.093 Å for the same symmetry. Natália et al. [49] performed classical MD and classical MC simulations and reported Ni-O bond lengths of 2.250 Å and 2.210 Å respectively. Aguilar et al.[51] reported the bond lengths of 2.05 Å and 2.092 Å from MC simulation and B3LYP/LANL2DZ calculation respectively. Mare et al. [53] performed calculations with various programs and basis sets, and obtained values of Ni-O as 2.073 Å (def2-SVP), 2.086 Å (def2-TZVP, def2-aug-TZVP, def2-QZVP), 2.089 Å (TZVP), 2.092 Å (QZV3P), 2.093 Å (TZV2P) and 2.100 Å (DZVP). Varadwaj et al. [55] reported a value of 2.093 Å by RX3LYP/6-311++G(d,p) calculation. Rulišek et al. [56] used B3LYP/6-311++G(d,p) to give a bond length of 2.087 Å. Akesson et al. [57] reported an average value of 2.108 Å at the HF-SCF level of theory. Fujii et al. [61] found a value of 2.086 Å from B3LYP/6-311+G(d,p) calculation. There are some experimental values of 2.002 Å [46], 2.060 Å [47], 2.10 Å [58] and 2.056 Å [48] determined in the liquid for the coordination of the Ni\(^{2+}\) ion with six water molecules. D’Angelo et al.[50] performed a XANES (X-ray absorption near-edge
structure) study on aquanickel(II) and reported a value of 2.03 Å. Inada et al. [52] reported a value of 2.05 Å from extended X-ray absorption experiment. Neilson et al. [54] found the Ni-O distance to be 2.06 Å by ND experiment. Waizumi et al. [62] reported the bond length to be 2.002 Å by XRD method. Magini et al. [63,64] and Caminiti et al. [65] also performed XRD experiments and found the value of 2.069 Å, 2.056 Å - 2.072 Å and 2.05 Å - 2.06 Å respectively in NiCl₂ solution. Sandstrom et al. [66] performed extended X-ray absorption fine structure study on NiCl₂ solutions and give a value of 2.07 Å for Ni-O distance. These computational and experimental data are all in good agreement with our results.

Vibrational stretching frequencies were also calculated for this set of hydrated nickel(II) complexes as seen in Figure 3-3 to 3-5. One of the reasons for calculating these complexes is so that they can be compared to experimental Raman studies that have been or will be completed. They can be used to confirm or refute any experimental Raman findings that cannot be unambiguously assigned to a specific complex. Solution Raman data that could be found regarding these complexes were for hexaaquanickel(II) ion. Fujii et al. [61] performed B3LYP/6-311+G(d,p) calculations and reported the totally symmetric vibrational frequencies to be 347 cm⁻¹, which is in consistent with our results of 335 cm⁻¹ (HF), 355 cm⁻¹ (MP2) and 350 cm⁻¹ (B3LYP). Edwards and Knowles [68] and Bickley et al. [69] performed Raman spectroscopy in formate solution and malonate solution respectively, and their results were 390 cm⁻¹ and 395 cm⁻¹. The calculated results underestimate the values obtained experimentally by 40 cm⁻¹- 60 cm⁻¹ due to the fact that our calculation did not include solvation effects. These underestimations are consistent
with similar studies completed previously on scandium chloro complexes [70] and zinc chloro complexes [71].

Some general trends can be drawn from the plots of bond lengths and stretching frequencies shown in Figures 3-3 to 3-5. Concerning the bond lengths, as more water ligands are added to the nickel, the Ni-O bond lengths show a general increasing trend. This is caused by overcrowding of the nickel resulting in the electronic repulsion between the water ligands. The opposite trend is seen in the Ni-O vibrational stretching frequencies. As the Ni-O bond lengths increase the vibrational stretching frequencies decrease. This trend is logical because if a bond is longer it will take less energy to cause it to vibrate and therefore have a corresponding frequency that is lower.
Figure 3-3: Ni-O bond lengths and vibrational stretching frequencies for [Ni(H$_2$O)$_6$]$^{2+}$, where $n$=1-6, calculated at the HF/6-31+G* level.
Figure 3-4: Ni-O bond lengths and vibrational stretching frequencies for [Ni(H$_2$O)$_n$]$^{2+}$, where n=1-6, calculated at the MP2/6-31+G* level.
Figure 3-5: Ni-O bond lengths and vibrational stretching frequencies for [Ni(H_2O)_n]^{2+}, where n=1-6, calculated at the B3LYP/6-31+G* level.
3.3 **Aquachloronickel(II) Complexes, \([\text{NiCl}_n(\text{H}_2\text{O})_m]^{2-n}\), where \(n=1-4\), \(m=0-(6-n)\)**

3.3.1 **Results**

Nickel chloride complexes are most frequently reported as neutral dichloronickel complex, which has been studied comprehensively by electron diffraction [73], ab initio methods [72,87], X-ray diffraction [74], crystal field study [88], Raman[90], Infrared[89], and UV-vis-NIR[78]. All these studies showed that the NiCl_2 molecule is linear. [85] Lemaire et al. [72] investigated NiCl_{n-2} complexes by DFT calculations, and suggested that the NiCl_3 is the most stable complex. The hydrated nickel chloride is known to form diaquadichloro- [74,91], tetraaquadichloro- [61,62,74,77,86], and pentaaquamonochloro-[61] species. Light green translucent crystals of NiCl_2(\text{H}_2\text{O})_4 (s) can be grown from an aqueous solution of NiCl_2. [77] In the neutral NiCl_2(\text{H}_2\text{O})_4 complex, the nickel(II) ion is surrounded by four water molecules and two chloride ions, and two of the water molecules are coordinated to nickel approximately in one of the tetrahedral directions; the other two are trigonally coordinated.[77] NiCl_2(\text{H}_2\text{O})_4 can be either the trans- or cis-isomer, where cis-NiCl_2(\text{H}_2\text{O})_4 is more stable than trans- NiCl_2(\text{H}_2\text{O})_4 without solvation, while trans-NiCl_2(\text{H}_2\text{O})_4 is more stable than cis-NiCl_2(\text{H}_2\text{O})_4 when solvation is included. [61] Thermodynamic studies of dehydration have been completed on dichloronickel(II) complexes. [91]

Stable structures were found for all anhydrous chloronickel(II) complexes up to and including tetra-coordinate species. These aquanickel complexes were hydrated and studied further up to and including hexa-coordinate species. The penta- and hexachloronickel were also studied but these structures underwent dissociation of one or
more chloride ligands. Therefore, stable structures for those complexes were not found. The total energies for the stable geometries, as well as the others that were attempted can be found in Table 3A.7 of the supplementary materials section. The geometries of the lowest energy structures can be found in Figure 3-6 for the mono- and dichloronickel complexes and Figure 3-7 for the tri- and tetrachloronickel complexes.

The anhydrous monochloronickel complex has $C_{\infty v}$ symmetry. The monohydrate has $C_{s}$ symmetry. The dihydrate has stable $C_{2v}$ structure at both levels, but the most stable structure is $C_{1}$. The most stable structure for the trihydrate is $C_{3}$ at MP2, but the same symmetry has an imaginary at B3LYP, and the preferred geometry is $C_{1}$ at B3LYP. The tetrahydrate has a preferred symmetry dependent on the level of theory. The most stable structure is $C_{2}$ for MP2, whereas the most stable structure at B3LYP is $C_{1}$. The most stable pentahydrate has no symmetry at both levels.
Figure 3-6: Optimized MP2 and B3LYP geometries for [NiCl₃(H₂O)ₓ]²⁻, where n=1-2 and m=0-(6-n). All symmetries marked with “*” indicate B3LYP otherwise all MP2 and B3LYP structure are similar. (grey=nickel, red=oxygen, white=hydrogen, green=chloride)
Figure 3-7: Optimized MP2 and B3LYP geometries for [NiCl$_n$(H$_2$O)$_m$]$^{2-n}$, where n=3-4 and m=0-(6-n).

The anhydrous dichloronickel complex has $D_{3h}$ symmetry. The stable structure of the monohydrate is $C_1$ at both levels. The dihydrate is stable with $C_{2v}$ symmetry. The stable structure of the trihydrate is without symmetry at $C_1$. The tetrahydrate has stable cis-$C_2$ and trans-$C_{4h}$ (with lowest energy) structures at both levels.
The anhydrous trichloronickel complex did not cause any symmetry problems, and at both levels showed a preferred $D_{3h}$ symmetry. The monohydrate has a stable $C_s$ structure at HF level, but it showed an imaginary frequency at MP2 and B3LYP levels. The $C_1$ symmetry was preferred at both levels. The stable structure of the dihydrate is without symmetry at $C_1$ at both levels. The trihydrate complex has lowest structure $C_{3v}$ at both levels. The meridional (mer) $C_1$ structure is also stable and competitive in energy.

The anhydrous tetrachloronickel complex did not cause any symmetry problems and at both levels showed a preferred $T_d$ symmetry. The stable structure of the monohydrate has $C_{2v}$ geometry with two chlorides hydrogen bonded to the hydrogen atoms respectively. The stable structure of the dihydrate has $C_{2v}$ symmetry with two chlorides hydrogen bonded with two hydrogen atoms from two different water molecules.

No stable structures were found for the penta- and hexachloronickel(II) complexes. All the geometries attempted showed dissociation of chloride ligands. Since stable structures were not found for these complexes, they will not be discussed further.

Plots of the Ni-O and Ni-Cl bond lengths and vibrational stretching frequencies were constructed and can be seen in Figure 3-8, Figure 3-9, and Figure 3-10 for HF, MP2 and B3LYP respectively. Within the plots, the lines highlighted in red indicate which bond lengths and stretching vibrations involve chloride ligands. Simulated polarized Raman spectra were also created for the stable HF structures and can be found in Figure 3A-2 of the supplementary materials section.
3.3.2 Discussion/Literature Comparison

Structural comparisons made between our results and the literature found on these complexes will be done primarily through bond length data. Lemaire et al. [72] reported from BLYP/6-31G* calculations of NiCl$_{n}^{2-n}$ (n=1-4) with the same symmetry as our results. The bond lengths reported by Lemaire et al. [72] are 2.11 Å, 2.11 Å, 2.23 Å, and 2.38 Å respectively (NiCl$_{n}^{2-n}$, n=1-4). Our results for NiCl$_3^-$ and NiCl$_4^{2-}$ are 2.24 Å and 2.40 Å at B3LYP level which are in good agreement with their results, and our results for NiCl$_2^+$ and NiCl$_2$ are 2.02 Å and 2.07 Å which are lower than their results about 0.09 Å – 0.04 Å at B3LYP level. Their study also showed that nickel complexes with more than four chlorine atoms are not stable. Hargittai et al. [73] reported the NiCl$_2$ bond distance of 2.076 Å by a combined gas-phase electron diffraction and vibrational spectroscopic experiment which is in consistent with our result. Ashworth et al. [75] reported the experimental value of NiCl$_2$ bond length to be 2.056 Å. Marcus [76] reported the Ni-Cl distance of NiCl$_2$ to be 2.07 Å. Waizumi et al. [74] reported the Ni-Cl bond length of NiCl$_2$ to be 2.06 Å. All these data are in good agreement with our results.

A few data have been found for hydrated nickel chloride. Fujii et al. [61] reported the Ni-O and Ni-Cl bond lengths for NiCl(H$_2$O)$_5^+$ to be 2.11-2.13 Å and 2.29 Å respectively by B3LYP/6-311+G(d,p) calculations. Fujii et al. [61] also calculated the Ni-O and Ni-Cl bond lengths for cis-NiCl$_2$(H$_2$O)$_4$ to be 2.108-2.109 Å and 2.347 Å respectively. They also reported the geometries for trans-NiCl$_2$(H$_2$O)$_4$, which has higher energy than the cis-isomer. Waizumi et al. [62] and Ptasiewicz-Bak et al. [77] performed X-ray diffraction experiments on NiCl$_2$(H$_2$O)$_4$ crystals, and obtained the Ni-O bond
lengths to be 2.063 Å and 2.089 Å, and the Ni-Cl bond lengths to be 2.375 Å and 2.400 Å respectively. Waizumi et al. [74] reported the Ni-O and Ni-Cl bond lengths of NiCl(H$_2$O)$_5^+$ to be 2.05 Å and 2.36 Å respectively by X-ray diffraction experiments. For the NiCl(H$_2$O)$_5^+$, our results are 2.0887-2.1088 Å (Ni-O) and 2.2507 Å (Ni-Cl) at MP2, and 2.1007-2.1366 Å (Ni-O) and 2.2904 Å (Ni-Cl) at B3LYP. For the cis-NiCl$_2$(H$_2$O)$_4$, our results are 2.082-2.135 Å (Ni-O) and 2.309 Å at MP2, and 2.104-2.152 Å (Ni-O) and 2.353 Å at B3LYP. Our results are in good agreement with the computational results, and underestimate the Ni-Cl bond lengths by 0.022-0.091 Å when compared with the XRD experimental results.

The plots of Ni-O and Ni-Cl bond lengths and vibrational stretching frequencies can be found in Figures 3-8 to 3-10, and some general trends can be seen from the plots. As more water molecules are added to the system, the Ni-O and Ni-Cl bond lengths increased and the vibrational stretching frequencies decreased respectively. The Ni-O bond lengths are uniformly shorter than the Ni-Cl bond length by 0.1-0.3 Å. As the Ni-O and Ni-Cl bond lengths increase, the bond becomes weaker and it requires less energy to stretch, so that the vibrational frequency becomes lower. The Raman active frequencies are important for comparison with experiment. Simulated polarized Raman plots are located in Figure 3A-2 in the supplementary materials section. These plots are based on Raman intensities calculated at the HF level of theory. Unfortunately no experimental Raman data was found to compare these results with.
Figure 3-8: Ni-O (red “+”) and Ni-Cl bond lengths and vibrational stretching frequencies for [NiCl$_n$(H$_2$O)$_m$]$^{2-n}$, where n=1-4, m=0-(6-n), calculated at the HF/6-31+G* level.
Figure 3-9: Ni-O (red “+”) and Ni-Cl bond lengths and vibrational stretching frequencies for \([\text{NiCl}_n(\text{H}_2\text{O})_m]^{2-n}\), where \(n=1-4\), \(m=0-(6-n)\), calculated at the MP2/6-31+G* level.
Figure 3-10: Ni-O (red “+”) and Ni-Cl bond lengths and vibrational stretching frequencies for [NiClₙ(H₂O)ₘ]²⁻, where n=1-4, m=0-(6-n), calculated at the B3LYP/6-31+G* level.
3.4 Aquahydroxynickel(II) Complexes, \([\text{Ni(OH)}_n(\text{H}_2\text{O})_m]^{2-n}\), where \(n=1-4\), \(m=0-(6-n)\)

Nickel radioisotopes are known to form poorly crystalline \(\text{Ni(OH)}_2\) under highly alkaline conditions (pH > 12.5). [97] Nickel hydroxide is widely used as the active material for positive electrode in alkaline secondary batteries. [99] Nickel hydroxide has a hexagonal layered structure with two polymorphs, \(\alpha\) and \(\beta\). [100] The \(\beta\)-\(\text{Ni(OH)}_2\) is well crystallized, and it possesses a brucite-like structure, whereas the \(\alpha\)-\(\text{Ni(OH)}_2\) is unstable, poorly crystallized and consists of a stacking of positively charged \(\text{Ni(OH)}_{2-x}\) layers, with intercalated anions and water molecules in the interlayer space to restore charge neutrality. [101] The \(\alpha\)-nickel hydroxide is a metastable phase and it changes rapidly to the \(\beta\)-form in strong alkali. [100] The nickel hydroxide has been studied comprehensively by experiments. Spectroscopic studies on nickel hydroxide have been completed by Infrared [94], UV/Vis spectroscopy [95], Extended X-ray absorption fine structure [96], diffuse reflectance spectroscopy (DRS) [97], X-ray diffraction [98], inductively coupled plasma atomic emission spectroscopy (ICP-AES) [99] and Raman spectroscopy[102,103]. There are additional reports on nickel hydroxide, like synthesis of the \(\alpha\)- [99,100] and \(\beta\)- [99,100,101] forms and solubility under various temperature and pH conditions [104,105]. The monohydroxide nickel(I) is used as an additive to increase the discharge capacity of Cd electrodes. [94] The \(\text{NiOH}^+\) has been studied by computational methods [60,92-94], and the bond lengths, binding energies and infrared spectra have been calculated. For \(\text{Ni(II)} \sim 10^{-8}\) M and pH ~8, higher order hydrolysis complexes like \(\text{Ni(OH)}_3^+, \text{Ni(OH)}_4^{2-}\) have been reported by Baes and Mesmer [106]. Solubility measurements of \(\text{NiO}\) have also shown the existence of tri- and tetrahydroxo nickel complexes, as well as the
existence of Ni₄(OH)₄⁺ in a high pH environment. [104] A computational study of the structures for [Ni(OH)(H₂O)₅]⁺ and [Ni(OH)₂(H₂O)₄] has been completed. [61] The structural, energetic and electronic properties of [Ni(OH)(H₂O)₅]⁺, [Ni(OH)₂(H₂O)₂], [Ni(OH)₂(H₂O)₃], [Ni(OH)₂(H₂O)₄] have also been studied by computational methods. [95]

3.4.1 Results

Stable structures were found for all hydroxide complexes up to and including the hexacoordinate tetrahydroxo species. Calculations were also completed on the penta- and hexahydroxo complexes, but these species showed dissociation of one or more hydroxide ligand and preferred to exist as ion pairs. Stable structures of these ion pair complexes were not found. Total energies, as well as C-PCM energies for all hydroxide complexes studied can be found in Table 3A-12 of the supplementary materials section. Optimized geometries of the stable complexes are located in Figure 3-11 for the mono- and dihydroxo complexes and in Figure 3-12 for the tri- and tetrahydroxo complexes.

The anhydrous monohydroxo complex did not cause any symmetry problems as all levels showed a preferred C₅ symmetry with the bent Ni-O-H angle. For the monohydrate, C₅ structure was tried first, and it is stable but higher in energy than the C₁ species at both levels. The C₁ structure looks similar to the C₅ structure, but the hydroxide ligand is slightly twisted. Two C₅ structures were tried for the dihydrate, but they both showed imaginary frequencies at both levels. The C₁ symmetry is preferred at both levels. Various C₅ structures were also tried for the trihydrate, and ultimately the only stable structure is C₁. The preferred tetrahydrate structure is C₁, and no stable structures with
higher symmetry were found. The most stable pentahydrate has no symmetry at both levels.

Figure 3-11: Optimized MP2 and B3LYP geometries for [Ni(OH)(H₂O)₅]⁺ and [Ni(OH)(H₂O)₃]. (grey=nickel, white=hydrogen, red=oxygen)
For the anhydrous dihydroxonickel, two $C_{2v}$ structures were tried first, but only the “W” shaped structure was stable at both levels, and a $C_2$ structure was proved to be an energy minimum. For the monohydrate, the $C_s$ structure was only a minimum at MP2 level, and the $C_1$ structure was preferred otherwise. The only stable geometry found for the dihydrate and trihydrate has $C_1$ symmetry at both levels. The tetrahydrate was stable with $C_1$ symmetry with a water molecule in the second hydration sphere. A cis-$C_1$ structure with all the ligands directly bonded to nickel is also stable but competitive in energy.

The anhydrous trihydroxonickel has stable $C_3$ and $C_1$ structures at both levels, and they are competitive in energy. The $C_3$ structure proved to be a minimum. When a single water adds to Ni(OH)$_3^-$, a $C_s$ structure is formed although it is not stable at both levels. The only stable structure found for the monohydrate has $C_1$ symmetry with a water molecule migrated to the second solvation shell. For the dihydrate, also the only stable structure found was $C_1$ with a water molecule in the second hydration sphere and hydrogen bonded to the other water molecule which is directly bonded to the central nickel and one of the hydroxide. For the trihydrate, a facial (fac) $C_{3v}$ structure with all the ligands directly bonded to the central nickel was located, but it is only stable at MP2 level. The three water molecules migrated to the second hydration sphere to form new $C_{3v}$ and $C_3$ (with lower energy) structure at B3LYP level.
For the anhydrous tetrahydroxonickel, stable structures found were $S_4$ and $C_2$, with $C_2$ being energetically preferred at both levels. The only stable structure found for monohydrate was also with $C_2$ symmetry in which the water molecule in the second hydration sphere and hydrogen bonded to the hydroxide ligands. The dihydrate was also...
only stable at C\textsubscript{2} geometry with two water molecules in the second hydrogen sphere and hydrogen bonded to the hydroxide ligands.

Bond length and vibrational stretching frequency data was also tabulated for all the stable aquahydroxo complexes and can be found in the supplementary materials section. Plots of the bond lengths and vibrational stretching frequencies can be found in Figure 3-13, Figure 3-14 and Figure 3-15 for HF, MP2 and B3LYP respectively. Simulated polarized Raman spectra can be found in Figure 3A-3 of the supplementary materials section.

3.4.2 Discussion/Literature Comparison

The geometries that have been found can primarily be confirmed by comparison with available data from crystal structures found in literatures, or other computational studies that have been completed on similar complexes. The Ni-O-H bond angels are always bent in all the aquahydroxo complexes at all levels.

Previous ab initio calculations can be compared to this work. Our results for the Ni-O bond length of Ni(OH)	extsuperscript{+} are 1.734 Å and 1.768 Å for the MP2 and B3LYP levels respectively. Magnusson and Moriarty [60] found 1.76 Å for the Ni-O distance which is in good agreement with our B3LYP calculations. Trachtman et al. [92] found that the Ni(OH)	extsuperscript{+} ion had a Ni-O distance of 1.779 Å which is slightly longer than our results. It could be explained by their relatively larger basis set (6-311++G**). Ricca and Bauschlicher [93] obtained the Ni-O bond length to be 1.712 Å which is about 0.022-0.056 Å as they used a different basis set ([8s4p3d] contraction of the (14s9p5d) primitive set). Wang et al. [94] reported the Ni-O distance to be 1.764 Å by B3LYP/6-
Although Wang et al. used a larger basis set, their result is consistent with our B3LYP result. No one has examined the partially hydrated structures, so only the pentahydrate structures can be compared. Fujii et al. [61] obtained 1.944 Å (Ni-OH) and 2.11-2.14 Å (Ni-OH₂) for the pentahydrate Ni(OH)⁺ by B3LYP/6-311+G calculations. Their results are in agreement with what been found for the Ni-O distances to be 1.926 Å (Ni-OH) and 2.090-2.119 Å (Ni-OH₂) at MP2 level, and 1.936 Å (Ni-OH) and 2.110-2.138 Å (Ni-OH₂) at B3LYP level respectively.

The calculated Ni-O bond distance for Ni(OH)₂ is 1.761 Å and 1.730 Å for MP2 and B3LYP respectively. Wang et al. [94] obtained a C₂h structure with the Ni-O distance of 1.737 Å using B3LYP/6-311++G(3df,3pd). Our C₁ structure is very similar to the C₂h structure, and the C₂h structure was also been tested, but there are imaginary frequencies at both levels. These results could also been compared to previous experimental data. Pandya et al. [96] examined the EXAFS data of β-Ni(OH)₂ and reported the Ni-O bond length to be 2.07 Å. Vespa et al. [97] found the Ni-O distance to be 2.06 Å by EXAFS. McEwen [98] found 2.14 Å for the Ni-O bond length from crystallographic analysis. All their results are much higher than our calculated data as our calculations are performed on single Ni(OH)₂ molecule in gas phase, and their experiments used nickel hydroxide crystals with the nickel surrounded by six equidistant oxygen atoms. Godelitsas et al. [95] reported a computational study containing results for the di-, tri, and tetrahydrate species using B3LYP/LANL2DZ calculations. The dihydrate dihydroxo complex that they found contained Ni-O bond lengths of 1.874 Å and 1.878 Å (Ni-OH) and 2.076 Å and 2.107 Å (Ni-OH₂). They also reported the Ni-O distances of the trihydrate to be 1.899-2.033 Å (Ni-OH) and 2.077-2.153 Å (Ni-OH₂). These results are in agreement with what been
found for MP2 as 1.860-1.863 Å (Ni-OH) and 2.104-2.134 Å (Ni-OH₂), B3LYP as 1.837-1.841 Å (Ni-OH) and 1.955-2.213 Å (Ni-OH₂) for the dihydrate; MP2 as 1.871-1.950 Å (Ni-OH) and 2.157-2.191 Å (Ni-OH₂), B3LYP as 1.842-1.953 Å (Ni-OH) and 2.176-2.238 Å for the trihydrate. The structure they reported for the tetrahydrate has two water molecules in the second hydration sphere, whereas our structure has only one water molecule in the second hydration shell. Fujii et al. [61] found 1.998 Å (Ni-OH) and 2.17-2.19 Å (Ni-OH₂) for the tetrahydrated nickel hydroxide which is slightly longer than our results of 1.924-1.957 Å (Ni-OH) and 2.102-2.155 Å (Ni-OH₂) at MP2; 1.933-1.968 Å (Ni-OH) and 2.104-2.179 Å (Ni-OH₂) at B3LYP due to the larger basis set (6-311+G(d,p)). There were no other literature reports found regarding the remaining nickel complexes which have been investigated in this study and therefore nothing could be compared with the optimized geometries.

Plots were constructed of the Ni-O bond lengths and vibrational stretching frequencies and can be found in Figures 3-13 to 3-15. Within each of these plots a trend can be seen of increasing Ni-O bond length and decreasing Ni-O vibrational stretching frequencies with addition of water ligands. The Ni-OH₂ and Ni-OH bond lengths separated clearly, and the Ni-OH bond lengths are always significantly shorter than that of Ni-OH₂ due to its stronger attraction to Ni²⁺, and the Ni-OH stretching frequencies are always about 200-400 cm⁻¹ larger than Ni-OH₂ stretching frequencies.

Raman spectroscopy will be the primary experimental technique used to identify these compounds at high temperatures and pressures, based on the Raman active bands that have been calculated. This experimental work has not been completed yet and no
Raman data was found in the literature related to this set of nickel(II) hydroxide complexes. However, predicted polarized Raman spectra have been constructed based on intensities calculated at the HF level of theory and can be found in Figure 3A-3 of the supplementary materials section.
Figure 3-13: Ni-O bond lengths and vibrational stretching frequencies for \([\text{Ni(OH)}_m(\text{H}_2\text{O})_n]^{2-m}\), where \(m=1-4, n=0-(6-n)\), calculated at the HF/6-31+G* level.
Figure 3-14: Ni-O bond lengths and vibrational stretching frequencies for $[\text{Ni(OH)}_m\text{(H}_2\text{O})_n]^{2-m}$, where $m=1-4$, $n=0-(6-n)$, calculated at the MP2/6-31+G* level.
Figure 3-15: Ni-O bond lengths and vibrational stretching frequencies for \([\text{Ni(OH)}_m(\text{H}_2\text{O})_n]^{2-m}\), where \(m=1-4\), \(n=0-(6-n)\), calculated at the B3LYP/6-31+G* level.
3.5 Amminenickel(II) Complexes, \([\text{Ni(NH}_3\text{)}_m\text{(H}_2\text{O})_n]^{2+}\), where \(m=1\text{-}6, n=0\text{-}(6-n)\)

A number of studies on amminenickel ions have been carried out by experimental or computational methods. Studies on crystal structures of Ni\(^{2+}\) complexes show that nickel(II) usually favors four-coordinate and six-coordinate complexes. [108] A solubility study of nickel(II) ammino complexation also showed that the mono- and diammine nickel(II) ions occur at free ammonia concentrations above 2 mmolkg\(^{-1}\) and 6.5 mmolkg\(^{-1}\) respectively. [105] The monoamine nickel(II) ion has been studied by computational methods [60, 67, 107], and the geometries, electronic structures, and binding energies have been reported. The hexaammine complex has been studied comprehensively by ab initio calculations and experiments. The experimental studies include \(^{14}\text{N} \text{NMR}\) and proton magnetic resonance to determine the exchange rate [111, 113], infrared spectroscopy [112, 114, 115], extended X-ray fine absorption structure [46], Raman spectroscopy [115, 116], and solubility and stability studies [105]. In addition, ab initio investigations involving DFT-RX3LYP calculations on low spin hexaammine complexes [55], DFT studies on ammonia adsorption on Ni\(_n\) clusters [107], and HF calculations on the reactivity of Ni\(^{2+}\) in solutions have been completed [110]. For the hydrated amminenickel(II) complexes, only the hexa-coordinated species are reported by computational methods in the literature. [55, 108, 110]. Varadwaj et al. [108] performed computational studies on \([\text{Ni(NH}_3\text{)}_n\text{(H}_2\text{O})_{6-n}]^{2+}\) complexes, and reported that the \([\text{Ni(H}_2\text{O})_6]^{2+}\) is predicted to be the least stable, and the successive replacement of H\(_2\)O by NH\(_3\) results in a monotonic increase in the stabilization of the complex. It is reported that the stability of nickel(II) ammonia complexes is generally greater than that of water complexes according to computational studies which can be explained by the diffuse
character of the NH₃ lone pair which enhance the dispersion and σ donation contributions.

[67]

3.5.1 Results

The nickel(II) ammine complexes were studied up to and including hexacoordinate species with and without water molecules. Stable geometries were found for all of the complexes studied. Total molecular energies as well as C-PCM energies of all the geometries studied can be found in Table 3A.16 of the supplementary materials section. The optimized geometries of the stable structures can be found in Figure 3-16 for the mono- and diammine complexes and in Figure 3-17 for the tri-, tetra-, penta- and hexaammine complexes.

The anhydrate monoamine posed no problems at the highest possible symmetry C₃ᵥ. For the monohydrate, two Cₛ structures were tried, but only the Cₛ species with the oxygen atoms symmetric to the plane was stable (at MP2/6-31+G*). A C₁ structure (looks like Cₛ, but the ammine is slightly twisted) is preferred at B3LYP level. Four Cₛ structures were tried for the dihydrate, and the one in which one water molecule in the O-Ni-N plane and the other bisected by the plane was stable (at MP2/6-31+G*), whereas it reduced to C₁ symmetry (although still looks similar to Cₛ) at B3LYP level. C₃ᵥ, C₃, Cₛ, and C₁ structures were sequentially attempted for the trihydrate, and only the C₁ structure is stable at both levels. For the tetrahydrate, a stable Cₛ structure was located, and the pentahydrate also has Cₛ symmetry.

The D₃d and D₃h geometries were tried for the anhydrate diammine complex, and the D₃h species is only stable at B3LYP level, while the D₃d structure has the lowest
energy at both levels. For the monohydrate, $C_{2v}$ structures were tried first, but they all have imaginary frequencies at both levels. After desymmetrization along the imaginary modes, the structure became stable at $C_2$ (B3LYP/6-31+G*) and the global minimum $C_1$ (both levels). For the dihydrate, the $C_{2v}$ structure was tried first, but it has imaginary modes at both levels. The $C_2$ symmetry is only stable at MP2, whereas $C_s$ structure is preferred at B3LYP level. The preferred trihydrate structure is $C_2$. For the tetrahydrate, high symmetry structures trans $C_{2h}$ and trans $C_{2v}$ were tried first, but the trans $C_{2h}$ structure is only stable at MP2 level, and trans $C_{2v}$ shows imaginary frequencies at both levels. Two trans $C_2$ structures were also attempted, and the one in which the rotation axis along with the O-Ni-O is stable at both levels, whereas the other in which the rotation axis only pass the Ni atom is only stable at MP2 level. A cis $C_s$ structure is energetically preferred at both levels. The $C_1$ structure is also stable and competitive in energy.

For the anhydrous triamine species, the $C_{3h}$ structure is only stable at MP2 level, and $C_3$ structure is energetically preferred at both levels. $C_s$ structure was tried for the monohydrate, but it showed imaginary frequencies at both levels. The $C_1$ structure (looks similar to $C_s$) is preferred at both levels. For the dihydrate, only the $C_s$ structure in which two ammines are symmetric with plane is stable at both levels. The highest possible symmetry $C_{3v}$ was tried for the trihydrate, but it has numerous imaginary frequencies at both levels. The fac $C_3$ structure proved to be a minimum. The mer $C_1$ structure is also stable at both levels and competitive in energy.
Figure 3-16: Optimized MP2 and B3LYP geometries for \([\text{Ni(NH}_3]_{m}\text{(H}_2\text{O})_n]\)^{2+}, where \(m=1-2, n=0-(6-m)\). “*” indicates MP2 and B3LYP structures with different symmetry are similar. (grey=nickel, red=oxygen, white=hydrogen, brown=nitrogen)
Figure 3-17: Optimized MP2 and B3LYP geometries for $[\text{Ni(NH}_3)_m(\text{H}_2\text{O})_n]^{2+}$, where $m=3-6$, $n=0-(6-m)$, “*” indicates MP2 and B3LYP structures with different symmetry are similar.
The anhydrous tetraammine complex has $C_{2v}$ symmetry. For the monohydrate, $C_{2v}$ and $C_2$ were tried and they are both unstable at both levels, whereas a $C_s$ structure is proved to be a minimum. The trans dihydrate is stable with $C_{2v}$ and $C_s$ symmetry at both levels, although $C_{2h}$ (has imaginary frequencies at both levels) and $C_2$ (revert back to $C_{2v}$ after optimization) were attempted, and a cis $C_s$ structure is energetically preferred at both levels.

For the anhydrous pentaammine species, $C_{3h}$, $C_{3v}$, $C_3$, $C_s$, and $C_1$ structures were sequentially tried, and only the $C_1$ and $C_s$ structures are stable at B3LYP level, whereas only the $C_1$ structure is stable at both levels. The $C_s$ structure is energetically preferred at B3LYP level. For the monohydrate, a $C_s$ structure is only stable at MP2 level with the lowest energy, and it reduced to $C_1$ symmetry at B3LYP level. Finally, the anhydrous hexaammine species has stable $C_1$ symmetry at both levels, and a $C_s$ structure only stable at B3LYP and is energetically preferred at this level.

Bond length and vibrational stretching frequency data were collected and plots were constructed which can be seen in Figure 3-18, Figure 3-19 and Figure 3-20 for HF, MP2 and B3LYP respectively. Raman intensities were calculated at the HF level and therefore simulated polarized Raman plots were able to be created and these can be found in the supplementary materials section in Figure 3A-4.

### 3.5.2 Discussion/Literature Comparison

The geometries that have been calculated for $\text{Ni(NH}_3)_n^{2+}$ (where $n=1-6$) can be compared to computational results reported in the literature. Pilme et al. [67] found 2.091 Å for the Ni-N distance in anhydrous monoamminenickel(II) species by ROB3LYP/6-
311+G(2d) calculations. Magnusson and Moriarty [60] reported the Ni-N bond length to be 2.081 Å based on MP2/6-311+G** calculations. Chen et al. [107] obtained 2.05 Å for the Ni-N distance in Ni(NH₃)²⁺ using the NRLMOL set of codes developed by Pederson et al. Our results for Ni-N distance are 2.005 Å (MP2) and 2.016 Å (B3LYP) which are slightly shorter due to our relatively smaller basis set. Chen et al. [107] also reported the Ni-N bond lengths for anhydrous di-, tri-, and tetraammine species which are 2.03 Å, 2.02-2.04 Å, and 2.06-2.19 Å respectively which are also slightly longer than our results of 1.979 Å, 2.014 Å and 2.060 Å for MP2, and 1.997 Å, 2.013 Å and 2.068-2.069 Å for B3LYP respectively. It should be noted that the structures which Chen et al. obtained did not use symmetry whereas our structures are constrained to the appropriate symmetry, so that their geometries are slightly different from ours. Aguilar et al. [51] obtained the Ni-N distance of Ni(NH₃)₆²⁺ to be 2.206 – 2.210 Å based on B3LYP/LANL2DZ calculations which are slightly longer than our results of 2.195 – 2.204 Å at B3LYP/6-31+G*, and our MP2 results (2.175-2.177 Å) are still lower. The basis set- LANL2DZ does not include diffuse or polarization functions. Varadwaj et al. [108] reported the Ni-N bond length to be 2.205 Å (mean value) by UX3LYP/6-311++G(d,p) calculations which is in good agreement with our result. [108] There were no literature reports found for anhydrous pentaammine complex.

For the aquaamminenickel(II) complexes, Varadwaj et al. [108] calculated the structures of [Ni(NH₃)₆(H₂O)₆₋₆]²⁺ (where n=0-6) at UX3LYP/6-311++G(d,p) level. Although they did not do the calculations with symmetry, they still obtained the similar geometries with ours, and their results are very comparable with our results. Results from
Varadwaj et al. [108] are tabulated in Table 3.1 along with the results from this research for comparison.

Table 3.1: Geometry comparison of hexacoordinate aquaammine complexes (B3LYP level) with results reported by Varadwaj et al. [108]. All bond lengths (Å) are averages where appropriate.

<table>
<thead>
<tr>
<th>Complex</th>
<th>Varadwaj et al. [108]</th>
<th>Our Results</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>Ni-N</td>
<td>Ni-O</td>
</tr>
<tr>
<td>Ni(NH$_3$)(H$_2$O)$_5^{2+}$</td>
<td>2.082</td>
<td>2.113</td>
</tr>
<tr>
<td>trans-Ni(NH$_3$)$_2$(H$_2$O)$_4^{2+}$</td>
<td>2.095</td>
<td>2.157</td>
</tr>
<tr>
<td>cis-Ni(NH$_3$)$_2$(H$_2$O)$_4^{2+}$</td>
<td>2.100</td>
<td>2.146</td>
</tr>
<tr>
<td>fac-Ni(NH$_3$)$_3$(H$_2$O)$_3^{2+}$</td>
<td>2.116</td>
<td>2.190</td>
</tr>
<tr>
<td>mer-Ni(NH$_3$)$_3$(H$_2$O)$_3^{2+}$</td>
<td>2.116</td>
<td>2.201</td>
</tr>
<tr>
<td>trans-Ni(NH$_3$)$_4$(H$_2$O)$_2^{2+}$</td>
<td>2.151</td>
<td>2.211</td>
</tr>
<tr>
<td>cis-Ni(NH$_3$)$_4$(H$_2$O)$_2^{2+}$</td>
<td>2.140</td>
<td>2.250</td>
</tr>
<tr>
<td>Ni(NH$_3$)$_5$(H$_2$O)$_3^{2+}$</td>
<td>2.169</td>
<td>2.323</td>
</tr>
</tbody>
</table>

Plots were created for the Ni-O and Ni-N bond lengths and vibrational stretching frequencies which can be seen in Figures 3-18 to 3-20. These plots showcase the trends in the bond lengths and vibrational stretching frequencies as more ligands, either ammonia or water, are added to the system. The Ni-N bond lengths are uniformly shorter than the Ni-O bond lengths that could be due to the fact that N has a lower electronegativity than O and is more polarizable, and hence NH$_3$ forms the stronger bonds to nickel. As more water molecules are added to the system, both of the Ni-O and Ni-N distances become longer. This could be explained by the mutual repulsion between water and ammonia attached to the same metal center will lead to increases in bond distances. The stretching frequencies show the opposite trends, the stretching frequencies become lower as more ligands are added to the system.
Schmidt et al. [115] reported the experimental Raman data for the Ni(NH$_3$)$_6^{2+}$ to be 370 cm$^{-1}$ for the Ni-N symmetric stretch. Our values are 312 cm$^{-1}$, 334 cm$^{-1}$, and 320 cm$^{-1}$ at HF, MP2 and B3LYP respectively which underestimate this by 35 - 58 cm$^{-1}$. It is known that there are systematic errors in calculated frequencies at HF, MP2 and B3LYP levels depending on the system and basis sets. [38] Pye et al. [70] also assumed that this underestimation is systematic for the mixed complexes. The polarized Raman data are also plotted which can be seen in Figure 3A-4.
Figure 3-18: Ni-O and Ni-N (red “+”) bond lengths and vibrational stretching frequencies for [Ni(NH$_3$)$_m$(H$_2$O)$_n$]$^{2+}$, where m=1-6, n=0-(6-n), calculated at the HF/6-31+G* level.
Figure 3-19: Ni-O and Ni-N (red “+”) bond lengths and vibrational stretching frequencies for [Ni(NH$_3$)$_m$(H$_2$O)$_n$]$^{2+}$, where $m=1-6$, $n=0-(6-n)$, calculated at the MP2/6-31+G* level.
Figure 3-20: Ni-O and Ni-N (red “+”) bond lengths and vibrational stretching frequencies for [Ni(NH$_3$)$_m$(H$_2$O)$_n$]$^{2+}$, where $m=1-6$, $n=0-(6-n)$, calculated at the B3LYP/6-31+G* level.
3.6 Chlorohydroxonickel(II) Complexes, [NiCl_m(OH)_n(H_2O)_l]^{2-m-n}, where m=1-3, n=1-(4-m), l=0-(6-m-n)

This set of complexes was studied up to and including the overall hexa-coordinate species with various combinations of chloride, hydroxide and water. All complexes have at least one chloride and one hydroxide ligand. Although it was attempted, no stable structures were found for species with a combination of chloride and hydroxide ligands greater than four. Many of these complexes tended to form ion pairs. No computational or experimental data was found in the literature therefore the discussion will be limited to comparison of the optimized structures as well as bond distance and vibrational stretching frequency trends.

3.6.1 Results

Total energies have been tabulated for all successfully optimized geometries and can be seen in Table 3A.5 in the supplementary materials section for this chapter. The optimized geometries of stable structures can be seen in Figure 3-21 for chlorohydroxo and chlorodihydroxo complexes, Figure 3-22 for chlorotrihydroxo, dichlorohydroxo, dichlorodihydroxo, dichlorotrihydroxo and trichlorohydroxo complexes.

The anhydrous chlorohydroxo complexes posed no problems at the highest possible symmetry C_s. Two C_s structures were tried for the monohydrate, but they both show imaginary frequencies at all levels, and C_1 geometry is preferred. The dihydrate has stable C_s and C_1 structures at both levels, and C_1 is also energetically preferred at both levels. For the trihydrate, two C_s structures were also tried first, but the planar one showed dissociation of chloride atom, whereas the other showed imaginary frequencies at
all levels. A $C_1$ structure with one water molecule in the second solvation shell and hydrogen bonded to chloride and the other water molecule is preferred at both levels. The tetrahydrate has trans $C_s$ (with lower energy) and cis $C_1$ symmetry with all the ligands directly bonded to the central nickel.

$C_{2v}$ structure was tried first for the anhydrous chlorodihydroxo complex, and it is unstable at all levels. The $C_2$ structure is stable at MP2 level but becomes $C_1$ (although looks similar to $C_2$) at B3LYP level. For the monohydrate, $C_{2v}$ and $C_2$ structures were attempted first, but neither of them were stable. The stable structures found are $C_s$ and $C_1$, and the $C_s$ structure is energetically preferred at both levels. For both $C_s$ and $C_1$ structures, all ligands are directly bonded to nickel at MP2 level, but one of the water molecules migrated to the second hydration sphere forming hydrogen bond with the other two water molecules at B3LYP level. For the dihydrate, $C_{2v}$ and $C_2$ structures were also tried first, and they both show imaginary frequencies at both levels, and $C_1$ structure is preferred at both levels. For the trihydrate, $C_{2v}$, $C_2$, $C_s$ and $C_1$ structures were sequentially attempted and only the $C_1$ structure with one water molecule in the second hydration sphere and hydrogen bonded to water and hydroxide.
The anhydrous chlorotrihydroxo complex has $C_3$ symmetry, although $C_{3v}$ was also tested and showed imaginary frequencies at all levels. Various $C_n$ structures were
attempted for the monohydrate, and all of them showed either dissociation or imaginary frequencies. A $C_1$ structure with one water molecule in the second hydration shell and hydrogen bonded to two hydroxides is stable at both levels. When one more water is added, a $C_1$ structure is formed with only all the hydroxides directly bonded to central nickel, two water molecules migrated to the second solvation sphere and hydrogen bonded to the same hydroxide, whereas the chloride in the third hydration shell and hydrogen bonded to the two water molecules.

Two $C_s$ geometries were tried for the anhydrous dichlorohydroxo complex. The planar one is only stable at B3LYP level, whereas the other is only stable at MP2 levels. A stable $C_1$ structure is located and is energetically preferred at both levels. Four $C_s$ structures were tried for the monohydrate, but they all show imaginary frequencies at both levels. The only stable structure been found is $C_1$. The dihydrate ($C_1$) and trihydrate ($C_1$) presented a similar result, except the trihydrate has the chloride in the second hydration sphere and hydrogen bonded to the three water molecules.

Two $C_{2v}$, two $C_2$, four $C_s$ and various $C_1$ structures were attempted for the dichlorodihydroxo complex. Both of the $C_{2v}$ structures and the $C_2$ structure with the Cl-Ni-Cl rotation axis showed dissociation of chloride ions. The $C_2$ structure in which the rotation axis only passes the nickel atom is stable at MP2 level. The only stable $C_s$ structure (MP2) is the one with the two hydroxides symmetric with the Cl-Ni-Cl plane and is energetically preferred at MP2 level. All the $C_1$ structures reverted back to $C_2$ or $C_s$ after optimization. For the B3LYP level, all the structures mentioned previously showed dissociation of chlorides or imaginary frequencies, so no stable structure was found at
this level. The stable structure for the monohydrate has $C_s$ symmetry with two chlorides in the second hydration shell and hydrogen bonded to the water and hydroxide. For the dihydrate, the highest possible symmetry $C_{2h}$ and $C_{2v}$ were tried first. The $C_{2h}$ structure showed dissociation of chlorides, whereas the $C_{2v}$ structure with two chlorides hydrogen bonded to water molecules is stable at both levels. 27 different $C_s$ structures were also attempted, and only the structures reverted back to $C_{2v}$ are stable whereas others all showed dissociations of chlorides or imaginary frequencies.

The trichloromonohydroxo complex is only stable with $C_1$ symmetry at both levels, as all the $C_s$ structures attempted are unstable. The monohydrate also has $C_1$ symmetry with the water molecule in the second solvation sphere and hydrogen bonded to hydroxide and chloride. For the dihydrate, various $C_s$ structures were attempted, and only the $C_s$ structure with two hydrogen bonded chlorides is stable at both levels.
Figure 3-22: Optimized MP2 and B3LYP geometries for [NiCl\(_m\)(OH)\(_n\)(H\(_2\)O)\(_l\)]\(^{2-m-n}\), where \(m=1-3\), \(n=0-(4-m)\), \(l=0-(6-m-n)\).
3.6.2 Discussion/Literature Comparison

No data were found in the literature to compare the stable geometries with. However comment on the trends observed within the bond length and vibrational stretching plots constructed at each of the three levels of theory could be given. Within the bond length figures there is a general trend of increasing Ni-O and Ni-Cl bond distance as more water molecules are added to the system. However, when there are ion pairs, the trend is reversed. This could be caused by the dissociated chloride anion hydrogen bonding to the central complex, thereby stabilizing the structure, resulting in shorter bond distances. These trends are found with all three levels of theory. As mentioned in previous sections, when the bond lengths increase, the vibrational stretching frequencies associated with those bonds decrease. As the bond length trend reverses, the vibrational stretching frequency trend also changes. This trend with respect to the frequencies is also found with all levels of theory.

Since the results are going to be primarily used to compare with experimental Raman data generated by the colleagues at the University of Guelph, the Raman spectra based on the Raman intensities calculated at the HF/6-31+G* level of theory are constructed. These plots will be beneficial in the assignment of the experimental Raman bands. These plots are located in Figure 3A-5 in the supplementary material section.
Figure 3-23: Ni-Cl and Ni-O (red “+”) bond lengths and vibrational stretching frequencies for [NiClₘ(OH)ₙ(H₂O)ᵢ]²⁻, where m=1-3, n=1-(4-n), and i=0-(6-m-n), calculated at HF/6-31+G* level.
Figure 3-24: Ni-Cl and Ni-O (red “+”) bond lengths and vibrational stretching frequencies for $[\text{NiCl}_m(\text{OH})_n(\text{H}_2\text{O})_l]^{2-m-n}$, where $m=1-3$, $n=1-(4-n)$, and $l=0-(6-m-n)$, calculated at MP2/6-31+G* level.
Figure 3-25: Ni-Cl and Ni-O (red “+”) bond lengths and vibrational stretching frequencies for \([\text{NiCl}_m(\text{OH})_n(\text{H}_2\text{O})_l]^{2-m-n}\), where \(m=1-3\), \(n=1-(4-n)\), and \(l=0-(6-m-n)\), calculated at B3LYP/6-31+G* level.
3.7 Chloroamminenickel(II) Complexes, $[\text{NiCl}_m(\text{NH}_3)_n(\text{H}_2\text{O})_l]^{2-n}$, where $m=1-4$, $n=1-(6-m)$, $l=0-(6-m-n)$

This set of complexes was studied up to and including the hexa-coordinate complexes with various combinations of chloride, ammonia and water ligands. Every complex has at least one chloride and one ammonia ligand. Stable structures were found for all of the other combinations, although a few of them were not stable as completely bound structures as they preferred to have hydrogen bonded water or ammonia ligands and in a couple cases an ion pair was preferred. Only a couple of experimental studies have been completed regarding bond length or vibrational stretching frequency data for this set of complexes. Therefore, the discussion of the chloroammine species will be restricted to the findings regarding geometries, bond length and vibrational stretching frequency trends.

3.7.1 Results

Total energies have been collected and tabulated for all optimized geometries attempted, and can be found in Table 3A.6 of the supplementary materials section. The optimized MP2 and B3LYP geometries for the stable structures can be seen in Figure 3-26 for all monochloroammine and dichloroammine complexes, Figure 3-27 for all trichloroammine and tetrachloroammine complexes.

The monochloromonoammine complex is stable at the highest possible symmetry $C_{3v}$ at both levels. For the monohydrate, four $C_s$ and $C_1$ structures were attempted. The $C_s$ structure in which the hydrogen atoms of water are in the symmetric plane with the hydrogen from ammine pointing to water is preferred at MP2 level, and $C_1$ structure is
preferred at B3LYP. For the dihydrate, various C\(_s\) structures were tried, but the only stable structure found is C\(_1\) at both levels. For the trihydrate, high symmetry structures like C\(_{3v}\) and C\(_3\) were tried first, but they all showed imaginary frequencies at both levels. A C\(_s\) structure is preferred at MP2 level, and it reduced to C\(_1\) (with the ammonia slightly twisted) at B3LYP level. The tetrahydride prefers C\(_1\) (trans and cis) symmetry as all the C\(_s\) structures tested showed imaginary frequencies. The trans C\(_1\) structure has lower energy.

For the chlorodiammine complex, C\(_{2v}\) and C\(_2\) were tried first. The C\(_{2v}\) structure showed imaginary frequencies at both levels, whereas C\(_2\) is only stable at B3LYP level. A C\(_s\) structure is stable at MP2 level, but C\(_1\) is energetically preferred at both levels. The monohydrate is only stable with C\(_1\) symmetry as all the structures with higher symmetry attempted are unstable. For the dihydrate, two C\(_{2v}\) structures were attempted, but neither of them is stable. A C\(_2\) structure is proved to be a minimum. Various C\(_s\) structures were tried for the trihydrate, and only a fac C\(_s\) structure is stable at both levels. The most stable structure found for the trihydrate is mer C\(_1\).

For the chlorotriammine complex, C\(_{3v}\) and C\(_3\) structures are unstable at both levels. Desymmetrization along the imaginary modes resulted in a stable C\(_s\) structure and is energetically preferred at MP2, whereas C\(_1\) geometry with slightly lower energy is preferred at B3LYP level. The monohydrate has C\(_1\) symmetry as all the C\(_s\) structures tried have imaginary frequencies at both levels. For the dihydrate, a mer C\(_s\) structure is only stable at B3LYP level, and mer C\(_1\) is stable at both levels. A fac C\(_1\) structure is energetically preferred.
The anhydrous chlorotetraammine complex has $C_4$ symmetry. For the monohydrate, trans $C_2$ is only stable at MP2 level, and trans $C_1$ is stable at both levels. A cis $C_1$ structure has the lowest energy. For the anhydrous chloropentaammine complex, various $C_s$ structures were tested, and only one of them is stable at B3LYP level. A $C_1$ structure is preferred at MP2 level.

For the anhydrous dichloroammine complex, two $C_s$ structures were attempted. The one with Cl-Ni-Cl symmetry plane is only stable at MP2, whereas the other is only stable at B3LYP level. For the monohydrate, various $C_s$ structures were tried, and only the one in which the chlorides are symmetric about the O-Ni-N plane is stable at both levels. The dihydrate also has $C_s$ symmetry in which nickel, oxygen atoms and nitrogen are in the symmetric plane. The trihydrate is only stable with fac $C_1$ and mer $C_1$ (with lower energy) symmetry.

For the dichlorodiammine complex, $C_{2v}$ structure is only stable at MP2 level, and $C_1$ is preferred at B3LYP. $C_{2v}$, $C_2$ and various $C_s$ structures has been tested for the monohydrate, but they all showed imaginary frequencies at both levels, and ultimately it has $C_1$ symmetry. For the dihydrate, two all-trans $C_{2h}$ structures were tried first, and the one in which the oxygen of water molecules are bisected by the horizontal plane is stable at MP2 level. An all-trans $C_1$ structure is stable at both levels. Cl-trans $C_1$ structure is energetically preferred at both levels, and NH$_3$-trans $C_2$, all-cis $C_1$ and H$_2$O-trans $C_{2v}$ structures are also stable and competitive in energy.

For the dichlorotriammine complex, high symmetry structures like $C_{3h}$, $C_{3v}$ and $C_3$ were attempted first, but none of them is stable. A $C_s$ structure is stable at MP2 level,
and it reduced to C₁ at B3LYP. The mer monohydrate has C₁ symmetry as all the C₃ structures tried are unstable. A fac C₃ structure has lower energy at B3LYP level, but it possesses an imaginary frequency at MP2 level, where a fac C₁ structure is preferred. For the anhydrous dichlorotetraammine complex, trans-C₄h and trans-C₂v are energetically preferred at B3LYP and MP2 respectively. Other stable structures as trans-C₂ (both levels), cis-C₂v (B3LYP), and cis-C₂ (both levels) are also found and competitive in energy.

Two C₃ᵥ structures were tried for trichloroammine complex, and only the one in which hydrogens are in the same direction with the chlorides is stable at B3LYP level. A stable C₃ structure is located and is energetically preferred at both levels. The monohydrate has C₁ symmetry as all the C₃ structures attempted are unstable. For the dihydrate, various C₃ structures were attempted, and only the one with a chloride in the second hydration sphere and hydrogen bonded to the water molecules is stable at both levels.

For the anhydrous trichlorodiammine complex, D₃h, C₃h, C₃, C₂v, C₂ and C₃ structures were attempted. The D₃h, C₃h and C₃ structures have imaginary frequencies at both levels. The C₂v structure in which the hydrogens are in the same directions with chlorides is only stable at B3LYP level, whereas the other is only stable at MP2 at is preferred at this level. C₂ structure reverted back to one of the C₂v structures after optimization, and C₃ structure is energetically preferred at B3LYP. The monohydrate has C₁ symmetry with one of the chloride in the second solvation shell and hydrogen bonded to ammonia and water. For the anhydrous trichlorotriammine complex, fac-C₃ᵥ structures
is only stable at B3LYP level, whereas the fac-C$_3$ structure is stable at both levels, and both of them have all the ligands directly bonded to central nickel. Attempts to create a mer isomer resulted in one of the chloride migrated to the second hydration sphere and hydrogen bonded to three ammonia ligands, forming a stable C$_1$ structures and is energetically preferred at both levels.

Figure 3-26: Optimized MP2 and B3LYP geometries for $[\text{NiCl}_m(\text{NH}_3)_n(\text{H}_2\text{O})_l]^{2-n}$, where $m=1-2$, $n=0-(6-m)$, $l=0-(6-m-n)$. All symmetries marked with "*" indicate B3LYP otherwise all MP2 and B3LYP structure are similar. (grey=nickel, red=oxygen, white=hydrogen, green=chloride)
Figure 3-26: (continued)
Figure 3-26: (continued)
Figure 3-26: (continued)
For the anhydrous tetrachloroammine complex, two $C_{3v}$ structures were tried first. The one in which the hydrogen atoms pointing to the opposite direction with the chlorides is only stable at MP2, whereas the other is only stable at B3LYP level. A $C_3$ structure is also only stable at B3LYP level and competitive in energy. A $C_s$ structure is preferred at MP2 level. The monohydrate has $C_1$ symmetry with water and one of the chlorides in the second solvation shell. The anhydrous tetrachlorodiammine complex is stable with $C_{2v}$ geometry in which one of the chlorides is hydrogen bonded to the ammonia ligands.
Figure 3-27: Optimized MP2 and B3LYP geometries for [NiCl\(_m\)(NH\(_3\))\(_n\)(H\(_2\)O)\(_l\)]\(^{2-n}\), where \(m=3-4\), \(n=0-(6-m)\), \(l=0-(6-m-n)\)
3.7.2 Discussion/Literature Comparison

To confirm the geometries that have been investigated, the best way is to compare with computational or experimental studies found in the literature. Stromberg et al. [117] performed B3LYP/BSI (double $\zeta$ valence quality) calculation on anhydrous dichloroamminenickel(II) and reported the Ni-Cl and Ni-O bond lengths to be 2.19 Å and 1.90 Å respectively which is comparable with our results of 2.15 Å and 2.04 Å at B3LYP level. Leineweber et al. reported the crystal structure of dichloroamminenickel(II) [118] and dichlorodiammine [119] by XRD. The reported Ni-O and Ni-Cl distances for the dichloroammine complex are 1.96-1.97 Å and 2.375-2.468 Å, and for dichlorodiammine complex are 2.089-2.513 Å. Our results for the dichloroammine complex are 2.05 Å and 2.11 Å at MP2 level, and 2.06 Å and 2.24 Å at B3LYP which are shorter than their results due to our calculations were performed on single molecule in gas phase whereas their experiments are for solid layered crystals with more ligands around a single nickel. Unfortunately no other studies were found containing the geometries of these mixed complexes and therefore there is nothing to compare with any of the remaining geometries.

Plots containing Ni-O, Ni-N and Ni-Cl bond lengths and vibrational stretching frequencies (Figure 3-28, Figure 3-29 and Figure 3-30) have been created and are located below. Each of these has a trend of increasing bond length, for a given complex, when more water molecules are added to the system. The trends become less clear when there are more chloride ligands. With more chloride ligands there is a greater tendency for water and chloride ligand dissociation and subsequent hydrogen bonding to the central
molecule. The effect is usually a small increase in stabilization (decrease in total energy) of the central molecule, thereby resulting in slightly shorter bond lengths. The inverse relationship between bond distance and stretching frequency is observed. A lot of the stretches were mixes of two or more of the Ni-Cl, Ni-N and Ni-O stretches as suggested by the overlap of the characteristic symbols.

Once again, calculated Raman spectra were constructed based on the intensities calculated at the HF level. There is no Raman data found in literature, but these will be useful for future work to be completed by our colleagues.
Figure 3-28: Ni-O (+), Ni-N (x) and Ni-Cl (\(\cdot\)) bond lengths and vibrational stretching frequencies for \([\text{NiCl}_m(\text{NH}_3)_n(\text{H}_2\text{O})_l]^{2-n}\), where \(m=1-4\), \(n=1-(6-m)\), \(l=0-(6-m-n)\), calculated at HF/6-31+G* level.
Figure 3-29: Ni-O (+), Ni-N (x) and Ni-Cl (|) bond lengths and vibrational stretching frequencies for \([\text{NiCl}_m(\text{NH}_3)_n(\text{H}_2\text{O})_l]^{2-n}\), where \(m=1-4, n=1-(6-m), l=0-(6-m-n)\), calculated at MP2/6-31+G* level.
Figure 3-30: Ni-O (+), Ni-N (x) and Ni-Cl (\(\mid\)) bond lengths and vibrational stretching frequencies for \([\text{NiCl}_m(\text{NH}_3)_n(\text{H}_2\text{O})_l]^{2-n}\), where \(m=1-4\), \(n=1-(6-m)\), \(l=0-(6-m-n)\), calculated at B3LYP/6-31+G* level.
3.8 Hydroxoamminenickel(II) Complexes, \([\text{Ni(OH)}_m(\text{NH}_3)_n(\text{H}_2\text{O})_l]^{2-m}\), where \(m=1-4\), \(n=1-(6-m)\), \(l=0-(6-m-n)\)

The hydroxoammine complexes were studied up to and including all the combinations of hexa-coordinate complexes. Each molecule investigated has at least one hydroxide ligand and one ammonia ligand. Structures were found for all combinations of ligands with four or less hydroxide ligands. When there are more than four hydroxide ligands, there is significant ligand dissociation. Some of the complexes do not have a completely bound set of ligands as hydrogen bonding is sometimes more energetically preferred. Very few studies were found in the literature that contained hydroxoammine complexes. Therefore, the discussion will be restricted to what have been found regarding geometries, bond length trends and vibrational stretching frequency trends.

3.8.1 Results

The total molecular energies for all optimized structures were tabulated in Table 3A.7 of the supplementary materials. The stable MP2 and B3LYP structures are in Figure 3-31 for the monohydroxoammine and dihydroxoammine complexes, and in Figure 3-32 for the trihydroxoammine and tetrahydroxoammine complexes.

The most stable monohydroxoammine has \(C_s\) symmetry, although other stable \(C_s\) and \(C_1\) structures exist with slightly higher energy. The monohydrate has \(C_1\) symmetry as all the \(C_s\) structures tried are unstable. The dihydrate and trihydrate are similar as they all have \(C_1\) symmetry at both levels. The tetrahydrate has stable trans \(C_1\) (with lower energy) and cis \(C_1\) structures.
Various C₅ structures were tried for the anhydrous hydroxodiammine complex, but they all have imaginary frequencies at both levels. A C₁ structure is preferred. The monohydrate is only stable with C₁ symmetry. The dihydrate has a stable C₅ structure at MP2 level, but a C₁ structure is energetically preferred at both levels. The trihydrate has stable fac-C₁ (with lower energy) and mer-C₁ structures.

The anhydrous hydroxotriammine complex has stable C₅ structure at both levels, and C₁ is preferred at B3LYP level. For the monohydrate, the C₅ structure has imaginary mode A" at both levels. Using tight optimization convergence criterion still resulted in a low imaginary mode A" (around -13 cm⁻¹) at both levels. A C₁ structure (looks similar to C₅) is preferred. For the dihydrate, a stable mer C₅ and two mer C₁ structure has been located, but a fac C₅ structure with lower energy is preferred.

The anhydrous hydroxotetraammine complex has C₅ symmetry. For the monohydrate, all the C₅ structures attempted have imaginary frequencies at both levels, and cis C₁ structure is preferred. A trans C₁ structure is also stable and competitive in energy. For the anhydrous hydroxopentaammine complex, the C₅ structure is only stable at B3LYP level, and it reduced to C₁ at MP2.

The anhydrous dihydroxoaammine complex is stable with the highest possible symmetry C₅. For the monohydrate, the C₅ structure is only stable at B3LYP level, and C₁ structure is preferred at both levels. The dihydrate (C₁) and trihydrate (C₁) presented similar results, except the trihydrate has a water molecule in the second hydration sphere and hydrogen bonded to water and hydroxide.
The anhydrous dihydroxodiammine is stable with its highest possible symmetry $C_{2h}$, $C_{2v}$, $C_2$, $C_s$ and $C_1$ were sequentially tested for the monohydrate, but only the $C_1$ structure is stable at both levels with a water molecule in the second solvation shell and hydrogen bonded to one of the hydroxide. For the dihydrate, stable structures found were one $C_s$ structure, two $C_2$ structures, and four $C_1$ structures. One of the $C_2$ structure in which two water molecules hydrogen bonded with only hydroxide ligands is preferred at both levels.

The anhydrous dihydroxotriammine complex is only stable with $C_1$ geometry. The monohydrate also has stable $C_1$ symmetry (preferred at MP2) with the water molecule in the second hydration sphere and hydrogen bonded to the ammonia and both of the two hydroxides. A fac $C_s$ structure is energetically preferred at B3LYP level. The anhydrous dihydroxotetraammine complex has stable trans $C_2$, trans $C_1$ (energetically preferred at B3LYP), and cis $C_2$ (with lowest energy at MP2).

For the anhydrous trihydroxoammine complex, two $C_{3v}$ structures were tried first. They both possessed numerous imaginary frequencies at both levels, and so are $C_3$ and $C_s$ structures. A $C_1$ structure is preferred at both levels. The monohydrate and dihydrate are both only stable with $C_1$ symmetry, and in which all the water molecules migrated to the second sphere and hydrogen bonded to hydroxides.

$C_{3h}$, $C_{3v}$, $C_3$, $C_s$ and $C_1$ structures were sequentially attempted for the anhydrous trihydroxodiammine complex, and only $C_1$ structure with one of the ammonia in the second solvation shell and hydrogen bonded to hydroxides. For the monohydrate, the added water molecules stayed with one of the ammonia in the second hydration sphere.
and hydrogen bonded to hydroxides forming a $C_1$ structure. For the trihydroxotriammine complexes, stable structures found were $C_3 \ [3+3]$ (energetically preferred at B3LYP) and $C_1 \ [4+2]$ (with lowest energy at MP2).

Figure 3-31: Optimized MP2 and B3LYP geometries for $\left[ \text{Ni(OH)}_m(\text{NH}_3)_n(\text{H}_2\text{O})\right]^{2-n}$, where $m=1-2$, $n=0-(6-m)$, $l=0-(6-m-n)$. (grey=nickel, white=hydrogen, brown=nitrogen)
Figure 3-31: (continued)
Figure 3-31: (continued)
Figure 3-31: (continued)
Figure 3-32: Optimized MP2 and B3LYP geometries for $[\text{Ni(OH)}_m(\text{NH}_3)_n(\text{H}_2\text{O})_l]^{2-n}$, where $m=3-4$, $n=0-(6-m)$, $l=0-(6-m-n)$
The anhydrate tetrahydroxoammine complex is stable with $C_1$ symmetry with the ammonia in the second solvation shell and hydrogen bonded to hydroxides. For the monohydrate, only the four hydroxides are directly bonded to central nickel, whereas the water and ammonia are in the second hydration shell and hydrogen bonded to hydroxides. The anhydrate tetrahydroxodiammine complex has $C_2$ symmetry with both of the ammonia hydrogen bonded to the hydroxides.

3.8.1 Discussion/Literature Comparison

No structural information was found in the literature for the hydroxoammine complexes, therefore there is no data to compare with what have been found in this research. However, the trends within the bond length and vibrational frequency plots those need is discussed. Plots containing Ni-O and Ni-N bond lengths and vibrational stretching frequencies (Figure 3-33, Figure 3-34 and Figure 3-35) have been created and are located below. The general trend within all the bond distance plots is that the Ni-O and Ni-N bond distances increase as more water molecules are added to the given complex. There was no ion pair formation among this set of complexes (i.e. no dissociation of hydroxide ligands). An inverse relationship exists between the bond distances and the vibrational stretching frequencies associated with those bonds. This is evident when looking at both the bond distance and stretching frequency plots side by side. As the bond lengths are increasing the vibrational stretching frequencies are decreasing to lower wavenumbers.

Raman intensity plots were created from the data obtained at the HF/6-31+G* level of theory. Although no experimental data could be found related to Raman activity
for this set of complexes, it is still of importance to our colleagues to have this information for comparison with their experimental Raman work in the future. Plots were constructed for all the HF energy minimum structures. Having this information will help in either assigning or confirming the experimental Raman bands.
Figure 3-33: Ni-O (|) and Ni-N (+) bond lengths and vibrational stretching frequencies for [Ni(OH)$_m$(NH$_3$)$_n$(H$_2$O)$_l$]$^{2-n}$, where $m$=1-4, $n$=1-(6-$m$), $l$=0-(6-$m$-$n$), calculated at HF/6-31+G* level.
Figure 3-34: Ni-O (\( \ddagger \)) and Ni-N (+) bond lengths and vibrational stretching frequencies for \([\text{Ni(OH)}_m(\text{NH}_3)_n(\text{H}_2\text{O})_l]^{2-n}\), where \(m=1-4\), \(n=1-(6-m)\), \(l=0-(6-m-n)\), calculated at MP2/6-31+G* level.
Figure 3-35: Ni-O (\(\|\)) and Ni-N (\(\mathbf{+}\)) bond lengths and vibrational stretching frequencies for \([\text{Ni(OH)}_m(\text{NH}_3)_n(\text{H}_2\text{O})_l]^{2-n}\), where \(m=1-4\), \(n=1-(6-m)\), \(l=0-(6-m-n)\), calculated at B3LYP/6-31+G* level.
3.9 Chlorohydroxoamminenickel(II) Complexes, \([\text{NiCl}_m(\text{OH})_n(\text{NH}_3)_l(\text{H}_2\text{O})_k]^{2-}\), where \(m=1-3\), \(n=1-(4-m)\), \(l=0-(6-m-n)\), \(k=0-(6-m-n-l)\).

This set of complexes was studied up to and including the overall hexa-coordinate species with various combinations of chloride, hydroxide, ammonia and water. All complexes have at least one chloride, one hydroxide, and one ammonia ligand. Many of these complexes tended to form ion pairs. No stable structures were found for anhydrous trihydroxochlorodi ammonium complex as all the possible geometries attempted resulted in dissociation of chloride ligand. No experimental or computational studies were found in the literature therefore the discussion will be limited to our findings regarding geometries, bond length and vibrational stretching frequency trends.

3.9.1 Results

Total energies have been tabulated for all successfully optimized geometries and can be seen in Table 3A.8 in the supplementary materials section. The optimized geometries of stable structures can be seen in Figure 3-36 for chlorohydroxoammine complexes, and Figure 3-37 for dichlorohydroxoammine and trichlorohydroxoammine complexes.

For the anhydrous chlorohydroxoammine complex, four \(C_s\) structures were attempted and only two of them are stable at both levels, and one is energetically preferred at MP2 level, whereas the other is preferred at B3LYP level. The monohydrate and dihydrate are both stable with \(C_1\) symmetry as all the \(C_s\) structures attempted are all unstable. For the trihydrate, both fac and mer (with lower energy) isomers are stable with \(C_1\) symmetry.
The anhydrous chlorohydroxodiammine complex has $C_1$ symmetry. The monohydrate is stable with $C_1$ symmetry with all the ligands directly bonded to central nickel(II) at MP2 level, but the water molecule dissociated and hydrogen bonded to hydroxide ligand at B3LYP level. For the dihydrate, depending on the trans-cis isomers, five stable $C_1$ structures and a $C_s$ structure were located, and two of which has water molecules involved in the formation of hydrogen bonds, and one of them is energetically preferred at both levels.

The anhydrous chlorohydroxotriammine complex is only stable with $C_1$ symmetry. For the monohydrate, a fac $C_1$ structure with all the ligands directly bonded to nickel is most stable at MP2 level, and a $C_1$ structure with the water molecule in the second hydration sphere is energetically preferred at B3LYP level. Various $C_s$ structures were tried for chlorohydroxotetraammine complex, and only one of the trans $C_s$ is stable at B3LYP level, and a trans $C_1$ structure is stable at both levels. A cis $C_s$ structure is energetically preferred at both levels.

Stable structures as $C_s$ and $C_1$ were located for anhydrous chlorodihydroxoammine complex, and $C_s$ has lower energy at B3LYP level whereas $C_1$ is energetically preferred at MP2. The monohydrate has a stable $C_s$ structure at MP2 level, but $C_1$ structure with water molecule in the second solvation shell and hydrogen bonded to hydroxide ligands has lower energy at both levels. The dihydrate was similar except the $C_s$ structure is stable at both levels and the $C_1$ structure (looks similar to $C_s$) with the chloride ligand hydrogen bonded to hydroxide and ammonia ligands instead of water molecule.
The anhydrous chlorodihydroxodiammine complex has stable C\textsubscript{2} (MP2) and C\textsubscript{1} (both levels with lower energy) structures and both of them have the chloride ligand in the second hydration sphere and hydrogen bonded to the ammonia ligands. The monohydrate has C\textsubscript{1} symmetry in which the water molecule is not coordinated to the central Ni\textsuperscript{2+} but is involved in the formation of hydrogen bond with coordinated hydroxide ligand. The anhydrous chlorodihydroxotriammine complex also have C\textsubscript{1} symmetry with the chloride and one of the ammonia ligands involved in the formation of hydrogen bonds with hydroxide and ammonia ligands.

Various geometries as C\textsubscript{3v}, C\textsubscript{3}, C\textsubscript{s} and C\textsubscript{1} were tried for chlorotrihydroxoammine complex, but only one of the C\textsubscript{1} structure with ammonia ligand in the second hydration sphere and hydrogen bonded to the hydroxide ligands is stable at MP2 level. All the structures attempted show dissociation of either ammonia or chloride ligand at B3LYP level, so no stable structures were found at this level. For the monohydrate, all the C\textsubscript{s} and C\textsubscript{1} structures tried also showed dissociation of chloride or ammonia ligand at both levels, so no stable structure was found for this complex. Also, no stable structures were located for chlorotrihydroxodiammine complex.

The anhydrous dichlorohydroxoammine complex is only stable with C\textsubscript{1} symmetry. The monohydrate is stable with its highest symmetry C\textsubscript{s}. Attempt to add a second water result in the formation of hydrogen bonds of water with chloride ligands and a C\textsubscript{1} structure is formed.
Figure 3-36: Optimized MP2 and B3LYP geometries for [NiCl(OH)ₙ(NH₃)ₙ(H₂O)ₖ]¹⁻ⁿ, where n=1-3, l=1-(5-n), k=0-(5-n-l). (grey=nickel, brown=nitrogen, green=chloride, white=hydrogen)
Figure 3-36: (continued)
Figure 3-36: (continued)
The dichlorohydroxodiammine complex has $C_s$ symmetry. The monohydrate is only stable with $C_1$ symmetry with one of the chloride ligand in the second solvation shell forming hydrogen bond with water and ammonia ligands. The anhydrous dichlorohydroxotriammine complex has stable $C_s$ ($\text{MP2 with all ligands directly bonded to nickel}$) and $C_1$ (both levels with lower energy and with the chloride hydrogen bonded with ammonia ligands).

Various $C_s$ structures were attempted for dichlorodihydroxoammine complex, but only the one in which both of chloride ligands hydrogen bonded to hydroxide and ammonia is stable. The monohydrate has $C_1$ symmetry and also has both of the chloride ligands involved in the formation of hydrogen bonds with ammonia and water. The dichlorodihydroxodiammine complexes is stable with $C_2$ symmetry in which the two chloride ligands are involved in the formation of hydrogen bonds with ammonia ligands.

The trichlorohydroxoammina complex is only stable with $C_1$ symmetry with the ammonia ligand hydrogen bonded to hydroxide. Attempt to add a single water result in two chloride ligands dissociated and hydrogen bonded to water and ammonia respectively. The anhydrous trichlorohydroxodiammine complex has $C_s$ symmetry with two chloride ligands in the second hydration sphere and hydrogen bonded to hydroxide and ammonia ligands.
Figure 3-37: Optimized MP2 and B3LYP geometries for [NiClₘ(OH)ₙ(NH₃)ₙ(H₂O)ₖ]²⁻ᵐ⁻ⁿ, where m=2-3, n=1-(4-m), l=0-(6-m-n), k=0-(6-m-n-l).
3.9.2 Discussion/Literature Comparison

The structures that have been investigated in this research could be compared with those in the literature to see if the calculated geometries are viable, however, no studies on this set of complexes were found in the literature. Thus only the trends of the bond lengths and vibrational stretching frequencies for each group of complexes could be discussed. Generally there is a trend that with more water molecules been added, the Ni-O, Ni-N and Ni-Cl bond lengths increase and stretching frequencies decrease. However, water, chloride or ammonia dissociated and involved in the formation of hydrogen bonds result in the shorter bond distances between those ligands which still directly bonded to nickel(II) and the central nickel(II), and in these cases the trends become less clear.

In most cases, no experimental or prior computational data exist for nickel complexes with more than two ligand types. However, favourable comparisons can be made for nickel complexes with one or two ligand types. Therefore, it is assumed that the more complicated nickel complexes have the same relative errors as the simpler nickel complexes.

The plots of Raman spectra which can be found in supplementary materials section were also constructed based on the Raman intensities for future experimental comparison by our colleagues at the University of Guelph.
Figure 3-38: Ni-O (+), Ni-Cl (|) and Ni-N (x) bond lengths and vibrational stretching frequencies for 

\[ \text{[NiCl}_m\text{(OH)}_n\text{(NH}_3\text{)}_l\text{(H}_2\text{O)}_k]^{2-m-n}, \]

where \( m=1-3, n=1-(4-m), l=0-(6-m-n), k=0-(6-m-n-l) \), calculated at HF/6-31+G* level.
Figure 3-39: Ni-O (+), Ni-Cl (⊥) and Ni-N (x) bond lengths and vibrational stretching frequencies for $[\text{NiCl}_m(\text{OH})_n(\text{NH}_3)_l(\text{H}_2\text{O})_k]^{2-m-n}$, where $m=1-3$, $n=1-(4-m)$, $l=0-(6-m-n)$, $k=0-(6-m-n-l)$, calculated at MP2/6-31+G* level.
Figure 3-40: Ni-O (+), Ni-Cl (|) and Ni-N (x) bond lengths and vibrational stretching frequencies for $[\text{NiCl}_m(\text{OH})_n(\text{NH}_3)(\text{H}_2\text{O})_l]^{2m-n}$, where $m=1-3$, $n=1-(4-m)$, $l=0-(6-m-n)$, $k=0-(6-m-n-l)$, calculated at B3LYP/6-31+G* level.
Chapter 4: Conclusion and Future Work

4.1 Conclusions

The present study was performed on possible nickel(II) complexes with all the combinations of chloride, hydroxide, water and ammonia ligands which may exist in CANDU Gen IV Supercritical Water-Cooled Reactors as corrosion products by ab initio computational chemistry methods. The geometries, total energies (with and without solvation effects), and vibrational stretching frequencies were obtained for all the stable complexes based on our calculations. Among most of the molecules with neutral or negative charge, chloride, ammonia or water ligands tend to be involved in the formation of hydrogen bonds instead of directly bonded to central nickel(II) ion based on our calculations. Although the Raman experimental results from the colleagues at the University of Guelph are still needed to compare with in order to predict the most likely corrosion products (as this study is performed on all the possible nickel complexes at zero Kelvin in the gas phase), the lists could also been narrowed down to a certain extent based on the calculation results and comparison with experimental data. Therefore, each set of complexes will be gone through and a list of what might be the best candidates will be provided.

4.1.1 Aquanickel(II)

Stable structures were found for all of the species up to and including the hexa-coordinate complex. The calculated structures are in good agreements with structural data and other calculation results. As discussed in Chapter three, the most possible complex to
exist in SCWR would be the hexaaquanickel(II) which has been reported in experimental literature.

4.1.2 Chloronickel(II)

Nickel(II) complexes with water and chloride ligands (from river water for the coolant system) were found up to and including the hexacoordinate species, and those complexes with more than four chloride ligands showed dissociation of chloride ions. In order to form deposition, it is needed to keep electronically neutral, so that the neutral complexes NiCl$_2$ and with water ligands up to and including hexacoordinate species are preferred. The hydrated tetrachloronickel(II) complexes tend to formed ion pairs indicating relatively instability. Other complexes with positive or negative charges in this set are also possible to exist in SCWR, and they need other cations or anions to form deposits.

4.1.3 Hydroxonickel(II)

This set of complexes are similar as the chloronickel(II) species in this case. Therefore, the most possible complexes to exist in SCWR would be the neutral Ni(OH)$_2$ and with one, two or three molecules. The tetrahydratedihydroxonickel(II), Ni(OH)$_2$(H$_2$O)$_4$, has one of the water molecules involved in the formation of hydrogen bond which indicate instability according to the calculations. The hydrated trihydroxo- and tetrahydroxonickel(II) complexes also have water molecules in the second hydration sphere. Thus, for the nickel(II) complexes which may exist in SCWR environment, the suggested complexes could be the monohydroxonickel(II) and its hydrated complexes. However, based on the calculations one of the water molecules migrated to the second
hydration sphere and hydrogen bonded to hydroxide within the trihydrate chlorohydroxonickel(II) indicating the instability.

4.1.4 **Amminenickel(II)**

This set of complexes are similar with hydrated nickel(II) complexes as in which all the ligands are directly bonded to central nickel(II) and with the same charge (2+). Furthermore, other computational studies reported that in this set of complexes, it becomes more stable when the water molecules are replaced by ammonia ligands. [108] The most possible complex to exist in SCWR would be the hexaamminenickel(II) which has been experimentally studied comprehensively. The hexaamminenickel(II) could also combine with other anions such as Cl to form neutral complex which has been reported in literature. [112]

4.1.5 **Chlorohydroxonickel(II)**

This set of complexes contain both chloride and hydroxide ligands and have studied up to and including the hexacoordinate species. The charges of this set of complexes are range from 0 to (-2). Therefore, the most probable complexes would be neutral NiCl(OH) and with water molecules up to hexacoordinate. Those complexes with negative charges of this set of complexes tend to form ion pairs which also suggested that they are unstable.

4.1.6 **Chloroamminenickel(II)**

This set of complexes contains neutral ammonia ligands and chloride ligands with a single negative charge, and stable structures were found up to and including the
hexacoordinate species. The neutral complexes in this set would be NiCl₂ with different combinations of water and ammonia ligands up to hexacoordinate species. Based on the calculations and comparison with experimental studies, the most possible complexes to exist in SCWR would be diaquadichlorodiamminenickel(II), dichlorotetraamminenickel(II), and dichlorodiamminenickel(II). The dichloromonoamminenickel(II) and its hydrated complexes are also possible to be the candidates in SCWR. When more chloride ligands are added to the central nickel(II), they tend to form ion pairs which suggest that they may be less stable.

4.1.7 Hydroxoamminenickel(II)

This set of complexes are similar with chloroamminenickel(II) to certain degrees. The neutral complexes would be Ni(OH)₂ with different combinations of water and ammonia ligands. As the nickel(II) complexes tend to form tetrahedral and octahedral species, the possible complexes to exist in SCWR would be Ni(OH)₂(NH₃)₂, Ni(OH)₂(NH₃)₂(H₂O) (H₂O in the second hydration shell), Ni(OH)₂(NH₃)₂(H₂O)₂ (both of H₂O in the second hydration shell) and Ni(OH)₂(NH₃)₄. Nickel(II) complexes with more than two hydroxide ligands tend to have water or ammonia ligands involved in the formation of hydrogen bonds. The monohydroxidenickel complex with ammonia and water could also exist if combine with other anions.

4.1.8 Chlorohydroxoamminenickel(II)

The complexes in this set contain at least one chloride, one hydroxide and one ammine ligands with 0, -1, or -2 charges, like the aquachlorohydroxocomplexes. The neutral complexes would be NiCl(OH) combine with different number of ammonia and
water ligands. As more hydroxide or chloride ligands been added, they tend to form ion pairs which suggest instability.

### 4.2 Future Work

Ab initio calculations were carried out on nickel(II) complexes with chloride, hydroxide, ammonia and water in this work. As the SCWR construction materials also contain other transition metal alloys, ions such as iron and chromium may be leached into environment and form complexes with those anions and molecules as well. Therefore ab initio investigations on iron(II) and chromium(II) complexes could be included in future work.
Reference:


