Measurement of $\pi^0$ photoproduction on the proton at MAMI C


Measurements of pion photoproduction on both proton and quasifree-neutron targets have a very long history, started about 65 years ago. Using the first bremsstrahlung facilities, pioneering results for $\gamma p \rightarrow \pi^0 p$ [1], $\gamma p \rightarrow \pi^+ n$ [2], and $\gamma n \rightarrow \pi^- p$ [3] were obtained. Despite all the shortcomings of the first measurements (such as large normalization uncertainties, wide energy and angular binning, limited angular coverage, etc.), those data were crucial for the discovery of the first excited nucleon state, $\Delta(1232)/2^+$ [4]. Though present electromagnetic facilities in combination with modern experimental detectors have allowed a

I. INTRODUCTION

Measurements of pion photoproduction on both proton and quasifree-neutron targets have a very long history, started about 65 years ago. Using the first bremsstrahlung facilities,
significant improvement in the quantity and quality of pion-photoproduction data [5], full understanding of pion-photoproduction dynamics is far from established, even in the best investigated area, comprising the first, second, and third resonance regions. The properties of many nucleon states in this energy range, especially in the third resonance region [6], are still not well understood. Resolving the complex partial-wave structure of pion photoproduction with minimal model assumptions requires more precise data for differential cross sections, with fine energy binning and full angular coverage, in combination with measurements of various polarization observables.

First results for polarization degrees of freedom have been recently reported by the A2, CBELSA/TAPS, and GRAAL Collaborations. The beam asymmetry $\Sigma$ has been measured by GRAAL [7] and CBELSA/TAPS [8]. Data for the other two single polarization observables $T$ (target polarization) and $P$ (recoil-nucleon polarization measured as a double-polarization observable with transversely polarized target and linearly polarized beam) have been reported from CBELSA/TAPS [9]. First results for double-polarization asymmetries have also been measured at CBELSA/TAPS with a longitudinally polarized target (observable $G$ with linearly polarized beam [8] and observable $E$ with circularly polarized beam [10]) and with a transversely polarized target and linearly polarized beam (observable $H$ [9]). The A2 Collaboration at MAMI obtained results for the double-polarization asymmetry $C^*$ using a circularly polarized photon beam and a nucleon recoil polarimeter [11]. Further data sets are currently under analysis, so that one can expect new analyses which provide much tighter constraints on the $\gamma p \to \pi^0 p$ reaction in the nearest future.

This work contributes to these efforts by presenting a new high-statistics measurement of the $\gamma p \to \pi^0 p$ differential cross sections conducted by the A2 Collaboration for incident-photon energies, $E_\gamma$, from 218 MeV up to 1573 MeV [or center-of-mass (c.m.) energies $W = 1136$–1957 MeV]. The data are obtained with a fine binning in $E_\gamma$ ($\sim$4 MeV for all energies below $E_\gamma = 1120$ MeV) and 30 angular bins, covering the full range of the $\pi^0$ production angle. The data obtained above $E_\gamma = 1443$ MeV ($W = 1894$ MeV) have limited angular coverage.

A more detailed analysis of the present $\gamma p \to \pi^0 p$ differential cross sections, combined with results from the A2 Collaboration obtained by measuring polarization observables, is currently in progress and will be published separately.

II. EXPERIMENTAL SETUP

The reaction $\gamma p \to \pi^0 p$ was measured using the Crystal Ball (CB) [12] as a central calorimeter and TAPS [13,14] as a forward calorimeter. These detectors were installed at the energy-tagged bremsstrahlung-photon beam produced by the electron beam of the Mainz Microtron (MAMI) [15,16].

The CB detector is a sphere consisting of 672 optically isolated NaI(Tl) crystals, shaped as truncated triangular pyramids, which point toward the center of the sphere. The crystals are arranged in two hemispheres that cover 93% of 4$\pi$ sr, sitting outside a central spherical cavity with a radius of 25 cm, which is designed to hold the target and inner detectors.

In the A2 experiments at MAMI C, TAPS was initially arranged in a plane consisting of 384 BaF$_2$ counters of hexagonal cross section. It was installed 1.5 m downstream of the CB center and covered the full azimuthal range for polar angles from $1^\circ$ to $20^\circ$. In the present experiments, 18 BaF$_2$ crystals, covering polar angles from $1^\circ$ to $5^\circ$, were replaced with 72 PbWO$_4$ crystals. This allowed running with a much higher MAMI electron current, without decreasing the TAPS efficiency due to the very high count rate in the crystals near the photon-beam line. The energy resolution of the PbWO$_4$ crystals has not been understood well yet, and the use of their information was restricted in the present analysis, as described later. More details on the calorimeters and their resolutions are given in Ref. [17], and references therein.

The present measurements used 1557-MeV and 1604-MeV electron beams from the Mainz Microtron, MAMI C [16]. The data with the 1557-MeV beam were taken in the first half of 2013 (Run-I) and those with the 1604-MeV beam in the second half of 2014 (Run-II). Bremsstrahlung photons, produced by the beam electrons in a 10-μm Cu radiator and collimated by a 4-mm-diameter Pb collimator, were incident on a 10-cm-long liquid hydrogen (LH$_2$) target located in the center of the CB. In Run-I with the 1557-MeV electron beam, the energies of the incident photons were analyzed up to 1448 MeV by detecting the postbremsstrahlung electrons in the Glasgow tagged-photon spectrometer (Glasgow tagger) [18–20]. The uncertainty in the energy of the tagged photons is mainly determined by the number of tagger focal-plane detectors in combination with the energy of the MAMI electron beam used in the experiments. Increasing the MAMI energy results in increasing both the energy range covered by the spectrometer and the corresponding uncertainty in $E_\gamma$. For the MAMI energy of 1557 MeV and the Glasgow tagger, such an uncertainty was about $\pm$2 MeV. The systematic uncertainty in the absolute value of $E_\gamma$, which is dominated by the energy calibration of the tagger, was about 0.5 MeV [20]. More details on the tagger energy calibration and uncertainties in the energies can be found in Ref. [21].

In Run-II with the 1604-MeV electron beam, the energies of the incident photons were analyzed from 1426 MeV up to 1576 MeV by detecting the postbremsstrahlung electrons in the Mainz end-point tagger (EPT) with 47 focal-plane detectors. Since the Glasgow tagger was not designed to measure the high-energy tail of the bremsstrahlung spectrum, the EPT spectrometer was built to conduct $\eta'$ measurements by covering this low-energy range of postbremsstrahlung electrons. The uncertainty of the EPT in $E_\gamma$ due to the width of its focal-plane detectors was about $\pm$1.6 MeV, with a similar value (i.e., $\sim$1.6 MeV) in the systematic uncertainty in $E_\gamma$ due to the EPT energy calibration. The energy calibration of the EPT is based only on the simulation of electron tracing, using measured magnetic-field maps. The correctness of this calibration, as well as its uncertainty, was checked by measuring the position of the $\eta'$ threshold, $E_\gamma \approx 1447$ MeV.

Because the main goal of Run-I was to measure the $\gamma p \to \pi^0 p$ reaction from $E_\gamma = 218$ MeV up to 1448 MeV and with full coverage of the $\pi^0$ production angle, the trigger required
the total energy deposited in the CB to exceed \(\sim 120\) MeV. Run-II was mainly dedicated to studying \(\eta^0\) physics. For that, the experimental trigger required the CB energy to exceed \(\sim 540\) MeV, rejecting the detection of \(\pi^0\) production in the very forward and backward angles.

### III. DATA HANDLING

The reaction \(\gamma p \rightarrow \pi^0 p \rightarrow \gamma \gamma p\) was searched for in events reconstructed with two or three clusters detected in the CB and TAPS together. In the analysis of the data from Run-I, the information from the 72 PbWO\(_4\) crystals of TAPS was not used at all. In the analysis of the data from Run-II, in which all PbWO\(_4\) crystals demonstrated good performance, their information was used in the cluster reconstruction. The two-cluster events were analyzed assuming that the final-state proton was not detected. This typically happens when the outgoing proton is stopped in the material of the downstream beam tunnel of the CB, or the proton kinetic energy in the laboratory system is below the software cluster threshold, which was 12 MeV in the present analysis.

The selection of event candidates and the reconstruction of the reaction kinematics were based on the kinematic-fit technique. Details on the kinematic-fit parametrization of the detector information and resolutions are given in Ref. [17]. The kinematic-fit technique was also used for the offline calibration of the calorimeters. The resolution functions used in the kinematic fit were adjusted to result in proper stretch functions (or pulls) and probability distributions. All two- and three-cluster events that satisfied the \(\gamma p \rightarrow \pi^0 p \rightarrow \gamma \gamma p\) hypothesis with a probability greater than 1% were accepted for further analysis. The kinematic-fit output was used to reconstruct the four-momenta of the outgoing particles. The events from the data of Run-II with more than one cluster in the 72 PbWO\(_4\) crystals of TAPS were discarded from the analysis, for the central range of \(\pi^0\) production angles. The resulting \(\gamma p \rightarrow \pi^0 p\) differential cross sections, all selected events were divided into 30 equal-width \(\cos \theta\) bins, covering the full range from \(-1\) to 1, where \(\theta\) is the angle between the directions of the outgoing \(\pi^0\) and the incident photon in the c.m. frame. Also, the quantity of data collected for \(\gamma p \rightarrow \pi^0 p\) was sufficient to obtain statistically accurate differential cross sections for every tagger channel in the energy range from \(E_\gamma = 218\) MeV to 1120 MeV, providing \(\sim 4\) MeV binning in \(E_\gamma\). For higher energies, energy bins were combined in two or more tagger channels.

The typical experimental statistics and acceptance are illustrated in Figs. 1, 2, and 3 for different energies and data sets. Figure 1(a) shows the experimental \(\cos \theta\) distribution for \(E_\gamma = 318\) MeV (\(W = 1215\) MeV), which is obtained after subtracting the random and the empty-target backgrounds. Since this energy bin is close to the \(\Delta(1232)/2^+\) maximum, the statistical uncertainties in this distribution, based on \(8 \times 10^3\) events, are very small (\(\sim 0.2\%\)). Figure 1(b) shows the corresponding angular acceptance, which is greater than 80\% for the central range of \(\pi^0\) production angles. The resulting \(\gamma p \rightarrow \pi^0 p\) differential cross section for \(E_\gamma = 318\) MeV is depicted in Fig. 1(c). It is also compared to the prediction from the partial-wave analysis (PWA) by SAID [23] with its CM12 [24] solution (shown by a solid line). The experimental data points lie slightly above the phenomenological prediction, but their angular dependence is very close to it. Because all modern PWAs give very similar predictions for the \(\gamma p \rightarrow \pi^0 p\) differential cross section at the \(\Delta(1232)/2^+\) maximum, the present results in this region do not contradict any of them.

Figure 2(a) shows the experimental \(\cos \theta\) distribution for \(E_\gamma = 1202\) MeV (\(W = 1771\) MeV), which is near the third resonance region. The statistical uncertainties in this distribution, based on \(8.5 \times 10^3\) events, are larger but are still sufficient for a reliable analysis of the resulting differential cross section. Since the shape of the \(\gamma p \rightarrow \pi^0 p\) differential cross sections in this region changes rapidly with energy, the main goal of the present work was to measure it with the finest energy binning. Figure 2(b) shows that the angular acceptance drops for the very forward angles. At these energies, this occurs as many \(\pi^0\) mesons produced at forward angles have both their decay photons going into TAPS, and such events do not pass the trigger requirements. The resulting differential cross section, plotted in Fig. 2(c), demonstrates again a quite reasonable agreement with the SAID CM12 solution [24].
FIG. 1. \( \cos \theta \) distributions for \( \gamma p \rightarrow \pi^0 p \), with \( \theta \) being the angle between the directions of the outgoing \( \pi^0 \) and the incident photon in the c.m. frame, obtained for \( E_\gamma = 318 \text{ MeV} \) \( (W = 1215 \text{ MeV}) \): (a) experimental data after subtraction of the empty-target background and the random background; (b) angular acceptance obtained from the MC simulation of \( \gamma p \rightarrow \pi^0 p \); (c) results for the \( \gamma p \rightarrow \pi^0 p \) differential cross section compared to the prediction by SAID \[23\] with its CM12 \[24\] solution. The error bars on all data points represent statistical uncertainties only.

Figure 3 shows the \( \cos \theta \) distributions for \( E_\gamma = 1512 \text{ MeV} \) \( (W = 1928 \text{ MeV}) \) from the data taken in Run-II. Despite the large quantity of the experimental events \( (9 \times 10^5) \), a very low, and even zero, acceptance for the very forward and backward angles, caused by the trigger, did not allow the first two and the last two \( \cos \theta \) bins to be measured. However, the remaining 26 data points of the differential cross section lie close to the prediction from SAID CM12 \[24\].

IV. EXPERIMENTAL RESULTS AND THEIR DISCUSSION

The \( \gamma p \rightarrow \pi^0 p \) differential cross sections were obtained by taking into account the number of protons in the LH2 target and the photon-beam flux from the tagging facilities corrected for the fraction rejected by collimators. The overall systematic uncertainty due to the calculation of the detection efficiency and the photon-beam flux was estimated to be 4% for the data taken in Run-I and 5% for the data taken in Run-II.

Another source of systematic uncertainty is due to imperfections in reproducing the angular dependence of the \( \cos \theta \) acceptance. Such imperfections are easier to see in differential cross sections with very low statistical uncertainties and a smooth shape. In the present data, it can be seen, for example, on the top of the experimental spectrum shown in Fig. 1(c). The angular dependence of differential cross sections can also be smeared due to the limited resolution in the \( \pi^0 \) production angle, resulting in a larger population of angular bins with a lower cross section. Such an effect is, possibly, seen in Fig. 3(c) near \( \cos \theta = 0.5 \).

Possible magnitudes of systematic uncertainties in differential cross sections can be checked by comparing similar results obtained by analyzing different data. Figure 4 illustrates self-consistency of the results for the \( \gamma p \rightarrow \pi^0 p \) differential cross sections obtained by the A2 Collaboration in three different experiments. The results for the lowest energy, \( E_\gamma = 218 \text{ MeV} \), taken in Run-I are compared in Fig. 4(a) with similar results obtained in the A2 experiment (A2-2008) with the Glasgow tagger and MAMI energy of 855 MeV \[25\]. As seen, the agreement of the two independent measurements is fairly good for both the absolute values and the angular dependence. Figure 4(b) compares the results of Run-I and A2-2008 for an energy near the \( \Delta(1232)/2^+ \) maximum, showing a small discrepancy in the angular dependence. Continuing the comparison of the results from Run-I and A2-2008 at energies above the \( \Delta(1232)/2^+ \) peak makes little sense as, at those energies, the A2-2008 data have much larger systematic uncertainties.

FIG. 2. Same as Fig. 1, but for \( E_\gamma = 1202 \text{ MeV} \) \( (W = 1771 \text{ MeV}) \).
uncertainties in the photon flux and in the angular acceptance, affected by the trigger. The results for two the highest energies, $E_\gamma = 1434$ and 1443 MeV, taken in Run-I are compared in Figs. 4(c) and 4(d) with similar results from the present work obtained in Run-II. Despite a general agreement, there are some discrepancies between the two sets of results, which cannot be explained by their statistical uncertainties. In our opinion, such discrepancies could be caused by several factors. There was a significant difference in the trigger of the two experiments and in the energy resolution of the calorimeters, as the experiment with the EPT (Run-II) was made with a much higher MAMI electron current. A poorer energy resolution in Run-II resulted in a larger smearing of the $\cos \theta$ angular dependence. Another possible explanation is a slight mismatch in Run-II resulted in a larger smearing of the $\cos \theta$ angular dependence. Overall, the average systematic uncertainty in the differential cross sections due to the resolution in the c.m. $\theta$ and possible imperfections in the angular acceptance was estimated to be 2% for the results of Run-I. The corresponding systematic uncertainty for the results of Run-II is at least twice as high.

Comparison of the results for the $\gamma p \rightarrow \pi^0 p$ differential cross sections from the present work with the most recent published data from other collaborations [7,26–29] is illustrated in Fig. 5 for eight different energies. Previous results with either large angular bins or very restricted angular coverage are not included in the comparison. As seen in Fig. 5, data with full angular coverage existed only in the region below the third resonance region. The results of this work are in a general agreement with all the data shown, although those data are not always entirely consistent with each other. As also evident, the present data from Run-I have much superior statistical accuracy, combined with a finer angular and energy binning.

The total cross sections for $\gamma p \rightarrow \pi^0 p$ were determined by integrating the differential cross sections for the data obtained from Run-I, which have the full angular coverage. The resulting total cross sections are plotted in Fig. 6 as a function of the c.m. energy $W$ along with earlier A2 results obtained for energies up to $E_\gamma = 444$ MeV ($W = 1309$ MeV) [25] and results from Ref. [26], obtained for energies up to $E_\gamma = 790$ MeV ($W = 1537$ MeV), and from Ref. [28], obtained for energies above $E_\gamma = 300$ MeV ($W = 1201$ MeV). The latter data have no full angular coverage; they were obtained with extrapolating experimental differential cross sections into forward and backward angles by using the corresponding PWA results. For completeness, the data at higher energies are also compared in Fig. 6(b) to the SAID solution CM12 [24], shown by a magenta dashed line. As seen in Figs. 6(a) and 6(b),

![Graphs and diagrams](image-url)
the present results are in good agreement, within the total uncertainties (calculated by adding the statistical and the overall systematic uncertainty in quadrature), with the earlier A2 results and in general agreement with the most recent experimental data and SAID CM2 solution based on fitting previous data, even if they do not have the full angular coverage.

Comparison of the present results for the $\gamma p \rightarrow \pi^0 p$ differential cross sections with predictions of different partial-wave and coupled-channel analyses is illustrated for selected energies in Fig. 7. Three sets of predictions shown in this figure correspond to PWAs (namely, SAID CM12 [24], MAID2007 [30], and Bonn-Gatchina BG2014-02 [31]) based on fitting previous data. The fourth set of curves comes from SAID PR15 solution obtained after adding the present data points into the SAID fit. Compared to the SAID solution CM12, the new fit PR15 also involves the data added to obtain SAID DU13 solution [5] and other most recent polarization data from CBELSA/TAPS [9,10] and A2 at MAMI [11]. As seen, all three the previous PWAs are very close to each other and to the present data points in the energy region around the $\Delta(1232)3/2^+$ maximum. Differences between the PWAs themselves as well as between the PWAs and the data points...
increase, especially for the MAID2007 solution, at energies that are away from the $\Delta(1232)3/2^+$ peak. Then all the PWAs and the data points become closer again in the second and the third resonance regions. The largest difference between the present data and the three previous PWAs is observed for the MAID2007 solution; this could be explained by the fact that many of the most recent data sets were not available at the time when the MAID2007 solution was released.

Quantitatively, the agreement of the present data with PWAs can usually be estimated via $\chi^2$ values calculated from deviations of experimental data points from PWA solutions. The $\chi^2$ function used in the minimization procedure by SAID is given by formula [32]

$$\chi^2 = \sum_i \left( \frac{X_{i} - \Theta_{i}^{\text{exp}}}{\epsilon_i} \right)^2 + \left( \frac{X - 1}{\epsilon_X} \right)^2,$$

(1)
TABLE I. Average $\chi^2/\mathrm{dp}$ values (including information on the total $\chi^2$ from Eq. (1) and the number of data points used) for MAID2007 [30] (fitting data only below $W = 2$ GeV) and three SAID solutions: CM12 [24], DU13 [5], and PR15.

<table>
<thead>
<tr>
<th>$\chi^2/\mathrm{dp}$</th>
<th>MAID2007</th>
<th>SAID CM12</th>
<th>SAID DU13</th>
<th>SAID PR15</th>
</tr>
</thead>
<tbody>
<tr>
<td>present $\gamma p \rightarrow \pi^0 p$</td>
<td>99722/7978 = 12.5</td>
<td>24036/7978 = 3.0</td>
<td>28745/7978 = 3.6</td>
<td>9434/7978 = 1.2</td>
</tr>
<tr>
<td>previous $\gamma p \rightarrow \pi^0 n$</td>
<td>135195/51468 = 8.7</td>
<td>73211/17087 = 4.3</td>
<td>41500/17087 = 2.4</td>
<td>40793/17087 = 2.4</td>
</tr>
<tr>
<td>$\gamma p \rightarrow \pi^+ n$</td>
<td>171853/8092 = 21.2</td>
<td>23533/8592 = 2.6</td>
<td>19312/8592 = 2.2</td>
<td>17540/8592 = 2.0</td>
</tr>
<tr>
<td>$\gamma n \rightarrow \pi^- p$</td>
<td>25335/2806 = 9.0</td>
<td>53657/3162 = 17.0</td>
<td>7561/3162 = 2.4</td>
<td>6572/3162 = 2.1</td>
</tr>
<tr>
<td>$\gamma n \rightarrow \pi^0 n$</td>
<td>22568/364 = 62.0</td>
<td>979/364 = 2.7</td>
<td>1436/364 = 4.0</td>
<td>1187/364 = 3.3</td>
</tr>
</tbody>
</table>

where $\Theta_i^{\exp}$ is an individual experimental value of the differential cross section with its uncertainty $\epsilon_i$ (includes both the statistical and the individual systematic uncertainty discussed in the text above), $\Theta_i$ is the fit value calculated for the same energy and the production angle, and $X$ is a normalization parameter with its uncertainty $\epsilon_X$ that is determined by the overall systematic uncertainties of the data (see start of this section). The average $\chi^2/\mathrm{dp}$ values, along with the information on the total $\chi^2$ from Eq. (1) and the number of data points used (dp), are listed in Table I for MAID2007 [30] and three SAID solutions: CM12 [24], DU13 [5], and PR15. As seen, the numbers in Table I confirm that the average $\chi^2/\mathrm{dp}$ is worse for the older solution MAID2007 than for the more recent SAID solutions CM12 and DU13. The magnitudes of $\chi^2/\mathrm{dp}$ obtained for the new SAID solution PR15 indicate that the new $\gamma p \rightarrow \pi^0 p$ data are consistent with the existing $\gamma N \rightarrow \pi N$ data sets. The number of data points available for each photoproduction reaction, which are listed in Table I, shows that the present $\gamma p \rightarrow \pi^0 p$ data increase the existing $\gamma p \rightarrow \pi^0 p$ statistics by $\sim 47\%$ and the existing $\gamma N \rightarrow \pi N$ statistics by $\sim 27\%$.

The $\chi^2/\mathrm{dp}$ values, averaged within each energy bin, are plotted as a function of $E_{\gamma}$ in Fig. 8(a) for the previous SAID solution DU13 [5] and for the new SAID PR15 solution, obtained after adding the present data into the fit. This comparison demonstrates the energies at which the new data give the most significant impact. The results obtained for the normalization parameter $X$ from the DU13 and PR15 solutions are plotted as a function of $E_{\gamma}$ in Fig. 8(b).

Another useful presentation of the new $\gamma p \rightarrow \pi^0 p$ results is obtained by plotting them as an excitation function at a particular production angle $\theta$ of $\pi^0$. In Fig. 9, the differential cross sections are shown as a function of the c.m. energy $W$ for ten of 30 measured $\cos \theta$ bins. As seen, the peaks corresponding to the first, the second, and the third resonance region are clearly evident in the excitation function for each $\pi^0$ production angle. The experimental distributions are also compared to the three previous PWA solutions (SAID CM12 [24], MAID2007 [30], and BG2014-02 [31]) and the new SAID PR15 solution. For a better comparison in the energy range away from the $\Delta(1232)/2^+$ peak, all distributions are also plotted after rescaling by a factor of 5. The comparison demonstrates that the largest discrepancies are seen for the very forward and backward production angles.

The strongest dip in the $\gamma p \rightarrow \pi^0 p$ excitation function is observed at very backward $\pi^0$ angles at the c.m. energies slightly below $W = 1.4$ GeV. An example of the $\gamma p \rightarrow \pi^0 p$ differential cross section near this dip (namely, at $W = 1380$ MeV) is shown in Fig. 10(a). This effect, caused by the nearly complete cancellation between different contributions to the production cross section, is even better seen in the results of the new SAID PR15 solution, shown in Fig. 10(b) by a blue solid line for $\cos \theta = -1$ (or $\theta = 180^\circ$). It is interesting that a very similar cancellation in the very backward

![FIG. 8](image_url)
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FIG. 9. (Color online) Present results for the $\gamma p \rightarrow \pi^0 p$ excitation function as a function of the c.m. energy $W$, shown for ten of 30 measured $\cos \theta$ bins (where $\theta$ is the $\pi^0$ production angle), are compared to previous PWA solutions from SAID CM12 [24] (blue dashed line), MAID2007 [30] (magenta long-dashed line), and BG2014-02 [31] (green dash-dotted line) and to a new SAID PR15 solution (red solid line). The error bars on all data points (black points for Run-I and cyan points for Run-II) represent statistical uncertainties only. For a better comparison, all distributions are also plotted after rescaling by a factor of 5.

$\pi^0$ production angles was observed in the charge-exchange reaction $\pi^- p \rightarrow \pi^0 n$ at the same energy range. The SAID W108 solution [33] (rescaled by a factor $10^{-3}$) for $\pi^- p \rightarrow \pi^0 n$ at $\theta = 180^\circ$ is plotted in Fig. 10(b) by a red dashed line. In contrast, the $\gamma p \rightarrow \pi^+ n$ reaction, depicted for the SAID PR15 solution by a black dashed line, has no such feature. Unfortunately, the quality and the angular coverage of the existing $\gamma n \rightarrow \pi^0 n$ data do not allow a reliable prediction for very backward production $\pi^0$ angles in this energy range [34].

For the c.m. energies above $W = 1.4$ GeV, the present results for very backward $\pi^0$ angles are in good agreement with similar results ($\theta = 170^\circ$) published in Ref. [35], also showing a cusp structure at the $\eta$-photoproduction threshold ($W = 1.486$ GeV).

The statistical accuracy of the present results is sufficient for reliable fits of the $\gamma p \rightarrow \pi^0 p$ differential cross sections with series of Legendre polynomials up to a high order. Of course, such decompositions are much less reliable without full angular coverage of the $\pi^0$ production angle. Thus, the present results from Run-II should be used only in energy-dependent PWAs with additional model assumptions. For Run-I, we describe the $\gamma p \rightarrow \pi^0 p$ differential cross sections of
each energy bin in terms of a series of orthogonal Legendre polynomials \( P_j(\cos \theta) \):

\[
\frac{d\sigma}{d\Omega}(W, \cos \theta) = \sum_{j=0}^{j_{\text{max}}} A_j(W) P_j(\cos \theta),
\]

with the energy-dependent decomposition coefficients \( A_j(W) \).

From the phenomenological point of view, the number of terms required in the series, \( j_{\text{max}} \), is related to the angular momenta of essential partial-wave amplitudes or to the highest spin of resonances, existing in the studied energy range. An isolated resonance with spin \( J = (2\ell + 1)/2 \) can contribute to coefficients \( A_j \) only with even \( j \) up to \( j = 2\ell \) and, hypothetically, to other coefficients via interference with sizable background amplitudes. For example, a wide four-star resonance \( \Delta(2420) \) with \( J^P = 11/2^+ \) \cite{6}, although lying above the energy range of the present measurements, can affect \( A_j \) up to \( j = 10 \) by itself and even higher \( A_j \) via the interference with other states. On the other hand, the maximum order \( j_{\text{max}} \), up to which such a decomposition is meaningful, is limited by the data quality.

In this work, the data of Run-I were fitted up to order \( j_{\text{max}} = 10 \). It was observed that including Legendre polynomials beyond \( j_{\text{max}} = 8 \) in the fits could not significantly improve the reduced \( \chi^2 \) (i.e., divided by the number of degrees of freedom, \( \chi^2/\text{nfd} \)). The comparison of the reduced \( \chi^2 \) for the fits of the data from Run-I with \( j_{\text{max}} = 6, 8, \) and 10 is shown in Fig. 11, using only the statistical uncertainties. In the energy range around the \( \Delta(1232)3/2^+ \) region, which is below \( W = 1.3 \) GeV, fitting with \( j_{\text{max}} = 6 \) should, in principle, be sufficient for a fairly good description of the differential cross sections. However, the high statistical accuracy, which arises from the very large cross sections, results in the reduced \( \chi^2 \) values varying mostly between two and three. As also seen in Fig. 11, including higher-order coefficients in the fits does not improve the situation much. It is necessary to take into account angular-dependent systematic uncertainties, with values varying between 1% and 2%, in order to obtain reasonable \( \chi^2 \) values in the \( \Delta(1232)3/2^+ \) region, even in the fits with \( j_{\text{max}} = 6 \). Above \( W = 1.7 \) GeV, using \( j_{\text{max}} = 6 \) is not sufficient, and the data require the minimum order of \( j_{\text{max}} = 8 \), corresponding to \( G \)-wave amplitudes.

The results of the Legendre-polynomial fits for each coefficient \( A_j \) are depicted in Fig. 12, showing their energy dependence in unprecedented resolution. The \( A_j \) error bars represent uncertainties obtained from fits that took into account only the statistical uncertainties in the cross sections. As previously discussed, the effect from the angular-dependent systematic uncertainties should be considered, especially for the \( \Delta(1232)3/2^+ \) energy region. To estimate the impact from these uncertainties, the values of the differential cross section for each angular bin were randomly shifted within the uncertainty magnitude (2%), followed by a new Legendre-polynomial fit. Such a procedure was then repeated several times, and the average spread of the fit results for each \( A_j \) was considered as its systematic uncertainty caused by the angular-dependent systematic uncertainties in the differential cross sections. The systematic uncertainties in \( A_j \), obtained from the fits with \( j_{\text{max}} = 10 \), are depicted by a black histogram at the bottom of each plot in Fig. 12, clearly peaking at the

![FIG. 10. (Color online) (a) present results for the \( \gamma p \rightarrow \pi^0 p \) differential cross section (blue full circles) at \( E_\gamma = 546 \) MeV (\( W = 1380 \) MeV) fitted with Legendre polynomials (red dashed line); (b) the new SAID PR15 solution (blue solid line) for \( \pi^0 \) production angle \( \theta = 180^\circ \) (cos \( \theta = -1 \)) as a function of the c.m. energy \( W \), compared with a similar \( \pi^- p \rightarrow \pi^0 n \) prediction (rescaled by a factor \( 10^{-3} \)) from SAID WI08 \[33\] (red dashed line) and a \( \gamma p \rightarrow \pi^+ n \) prediction (rescaled by a factor 0.3) from SAID PR15 (black dashed line).](024617-10)
position of $\Delta(1232)/2^+$. In this region, the magnitudes of the systematic uncertainties in $A_j$ at high orders are very close to the magnitudes of $A_j$ themselves. The Legendre-polynomial fits for the c.m. energies below $W = 1280$ MeV were made for both $j_{\text{max}} = 6$ and $j_{\text{max}} = 10$. Their comparison showed good agreement of the results for $A_0$ to $A_6$.

One of the most obvious features in the energy dependence of $A_j$ is the absence of any well-defined sharp structures, which could be associated with unknown narrow resonances. The results for $A_0$ show good agreement with the values for the total cross sections obtained by the integration of the differential cross sections, providing another confirmation of the good quality of the present data. The $A_0$ energy dependence clearly demonstrates the first, second, and third resonance regions, but appears structureless at higher energies. In contrast, some of the other $A_j$ reveal structures even at the energies above the third resonance region. These structures could be associated with contributions from $N^*$ and/or $\Delta^*$ states, with spins up to $7/2$ or higher, that are known in this region or near it [6]. These resonances have masses close to the upper end of the studied energy range, as $\Delta(1950)/2^+$, or even above it, as $N(2190)/2^-$, $N(2220)/2^+$, and $N(2250)/2^-$, but their large widths ($\sim 300$ MeV or even larger) make those resonance contributions noticeable in the present data.

In Fig. 12, the coefficients $A_j$ are also compared to results obtained from the predictions generated with SAID CM12 [24], MAID2007 [30], BG2014-02 [31], and SAID PR15, by fitting them, after adding small uncertainties, with Legendre polynomials up to $j_{\text{max}} = 10$. As seen, the Legendre-polynomial fits to the predictions of all models reproduce fairly well only the coefficients $A_0$ and $A_2$. Partial agreement is observed for $A_1$ and the coefficients $A_3$ to $A_7$. The coefficients $A_8$ to $A_{10}$ are small and show some meaningful structures only above $W = 1.7$ GeV.

Of course, such a Legendre decomposition of the present differential cross sections is only a very first step towards a full single-energy partial-wave analysis. A direct extraction of partial-wave amplitudes by simultaneous fitting the present differential cross sections and polarization observables will be performed in the future. However, the present analysis already shows the sensitivity of the new data to partial-wave amplitudes up to $G$ waves ($\ell = 4$).
V. SUMMARY AND CONCLUSIONS

New measurements of $\pi^0$ photoproduction on the proton have been conducted with the A2 tagged-photon facilities at energies provided by the Mainz Microtron, MAMI C, with a maximum electron-beam energy of 1.6 GeV. The differential cross sections, obtained with a fine energy and angular binning, increase the existing quantity of $\pi^0$ photoproduction data by $\sim 47\%$. A new solution PR15 of the SAID partial-wave analysis obtained after adding the new data into the fit is presented. The magnitudes of $\chi^2/\nu$ obtained for this solution indicate that the new $\gamma p \rightarrow \pi^0 p$ data are consistent with the existing $\gamma N \rightarrow \pi N$ data sets. Owing to the unprecedented statistical accuracy and the full angular coverage, the results are sensitive to high partial-wave amplitudes. This is demonstrated by the decomposition of the differential cross sections in terms of Legendre polynomials and by further comparison to model predictions. The present data are expected to be invaluable for future partial-wave and coupled-channel analyses, which could provide much stronger constraints on the properties of nucleon states known in this energy range and, perhaps, even reveal new resonances. A more detailed analysis of the present $\gamma p \rightarrow \pi^0 p$ differential cross sections, combined with results from the A2 Collaboration obtained by measuring polarization observables, is currently in progress and will be published separately.

ACKNOWLEDGMENTS

We thank A. Donnachie for useful remarks and continuous interest in the paper. The authors wish to acknowledge the excellent support of the accelerator group and operators of MAMI. This work was supported by the Deutsche Forschungsgemeinschaft (SFB443, SFB/TR16, and SFB1044), the European Community-Research Infrastructure Activity under the FP6 “Structuring the European Research Area” programme (Hadron Physics, Contract No. RII3-CT-2004-506078), Schweizerischer Nationalfonds (Contract Nos. 200020-156983, 132799, 121781, 117601, 113511), the UK Science and Technology Facilities Council (STFC 57071/1, 50727/1), the US Department of Energy (Offices of Science and Nuclear Physics, Award Nos. DE-FG02-99-ER41110, DE-FG02-88ER40415, DE-FG02-01-ER41194), and National Science Foundation (Grant Nos. PHY-1039130, IIA-1358175, INFN (Italy), and NSERC (Canada). Y.I.A. acknowledges support by the Russian Science Foundation (Grant No. 14-22-00281). A.F. acknowledges additional support from the TPU (Grant No. LRU-FTI-123-2014) and the MSE Program “Nauka” (Project No. 3.825.2014/K). We thank the undergraduate students of Mount Allison University and The George Washington University for their assistance.